

# A shock-capturing discontinuous Galerkin method for the numerical solution of inviscid compressible flow

Jiří Hozman

Technical University of Liberec  
Faculty of Science, Humanities and Education

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# Outline

- 1 System of the Euler equations
- 2 Discretization of the problem
- 3 Shock-capturing techniques
- 4 Numerical examples

# Introduction

- **Our aim:** efficient, accurate and robust numerical scheme for the simulation of **inviscid** and/or **viscous compressible flows**,
- **Model problem:** system of the Euler equations,
- piecewise regular solution, usually contains discontinuities,
- discontinuous Galerkin method (**DGM**),
- piecewise polynomial discontinuous approximation,
- **shock-capturing** scheme  $\sim$  system of the Navier-Stokes equations,
- artificial viscosity limiter  $\sim$  **residual of the entropy equation**,
- DGM with IPG variants: nonsymmetric, symmetric and incomplete.

- 1 **System of the Euler equations**
- 2 **Discretization of the problem**
  - Triangulations
  - DGM spaces
  - Space semi-discretization
  - Full time-space discretization
- 3 **Shock-capturing techniques**
  - Artificial viscosity limiter
  - Entropy-based viscosity
  - Discretization of shock-capturing scheme
- 4 **Numerical examples**
  - Implementation aspects of BDF-DG method
  - Numerical example 1
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  - Numerical example 3

# Governing equations (1)

## Compressible Euler equations

$$\frac{\partial \mathbf{w}}{\partial t} + \sum_{s=1}^d \frac{\partial}{\partial x_s} \mathbf{f}_s(\mathbf{w}) = 0 \quad \text{in } \Omega \times (0, T) \quad (1)$$

- $\mathbf{w} = (\rho, \rho v_1, \dots, \rho v_d, e)^T \in \mathbf{R}^{d+2}$
- Euler fluxes  $\mathbf{f}_s : \mathbf{R}^{d+2} \rightarrow \mathbf{R}^{d+2}$ ,  $s = 1, \dots, d$

$$\mathbf{f}_s(\mathbf{w}) = (\rho v_s, \rho v_s v_1 + p \delta_{s1}, \dots, \rho v_s v_d + p \delta_{sd}, (e + p) v_s)^T$$

- properties of fluxes  $\mathbf{f}_s$ :

$$\mathbf{f}_s(\xi \mathbf{w}) = \xi \mathbf{f}_s(\mathbf{w}), \quad \xi \in \mathbf{R}, \quad \xi \neq 0,$$

$$\mathbf{f}_s(\mathbf{w}) = \mathbf{A}_s(\mathbf{w}) \mathbf{w}, \quad s = 1, \dots, d,$$

$$\mathbf{P} = \sum_{s=1}^d \mathbf{A}_s(\mathbf{w}) n_s = \mathbf{T} \mathbf{\Lambda} \mathbf{T}^{-1} = \mathbf{P}^+(\mathbf{w}, \vec{n}) + \mathbf{P}^-(\mathbf{w}, \vec{n}),$$

## Governing equations (2)

- state equation for perfect gas:

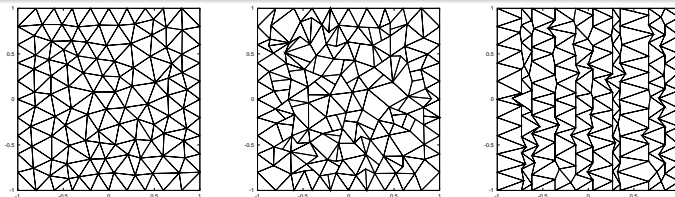
$$p = (\gamma - 1) (e - \rho |\mathbf{v}|^2 / 2)$$

- initial condition (IC):  $\mathbf{w}(x, 0) = \mathbf{w}^0(x)$  in  $\Omega$
- boundary conditions (BC):

boundary	character	extrapolated	prescribed
$\partial\Omega_i$ (inlet)	supersonic	—	$\rho, v_1, \dots, v_d, p$
	subsonic	$p$	$\rho, v_1, \dots, v_d$
$\partial\Omega_o$ (outlet)	supersonic	$\rho, v_1, \dots, v_d, p$	—
	subsonic	$\rho, v_1, \dots, v_d$	$p$
$\partial\Omega_w$ (walls)	$\mathbf{v} \cdot \vec{n} = 0$ (impermeability condition)		

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# Triangulations

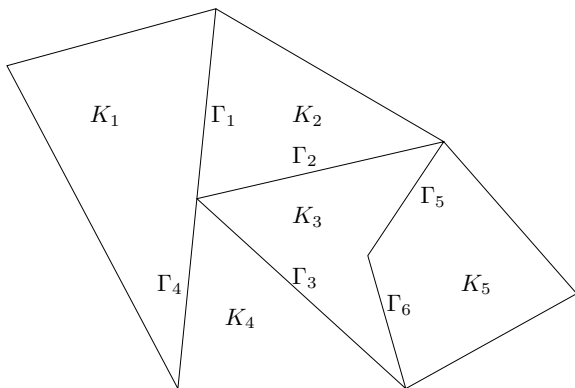


- let  $\mathcal{T}_h$ ,  $h > 0$  be a partition of  $\bar{\Omega}$
- $\mathcal{T}_h = \{K\}_{K \in \mathcal{T}_h}$ ,  $K$  are polygons/polyhedra
- let  $\mathcal{F}_h = \{\Gamma\}_{\Gamma \in \mathcal{F}_h}$  be a set of all edges/faces of  $\mathcal{T}_h$ ,
- we distinguish
  - **inner** edges/faces  $\mathcal{F}_h^I$ ,
  - **'Dirichlet'** edges/faces  $\mathcal{F}_h^D$  (i.e., inlet and solid walls),
  - **'Neumann'** edges/faces  $\mathcal{F}_h^N$  (i.e., outlet),
- we put  $\mathcal{F}_h^{ID} \equiv \mathcal{F}_h^I \cup \mathcal{F}_h^D$ .



# Fictional triangulation

- convex/nonconvex elements with/without hanging nodes



## Spaces of discontinuous functions

- let  $s_K \geq 1$ ,  $K \in \mathcal{T}_h$  denote local Sobolev index,
- let  $p_K \geq 1$ ,  $K \in \mathcal{T}_h$  be local polynomial degree,
- we set vectors  $\mathbf{s} \equiv \{s_K, K \in \mathcal{T}_h\}$  and  $\mathbf{p} \equiv \{p_K, K \in \mathcal{T}_h\}$
- over  $\mathcal{T}_h$  we define:
  - *broken Sobolev space*

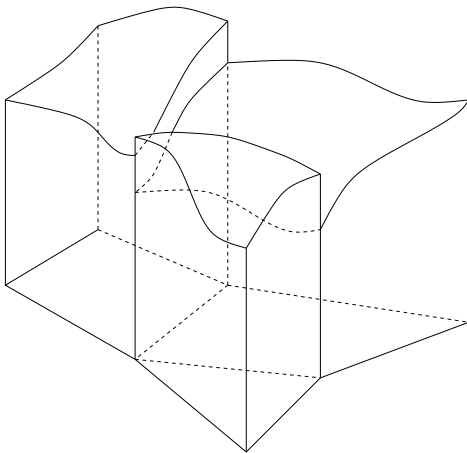
$$H^{\mathbf{s}}(\Omega, \mathcal{T}_h) = \{v; v|_K \in H^{s_K}(K) \forall K \in \mathcal{T}_h\}$$

- the space of piecewise polynomial functions

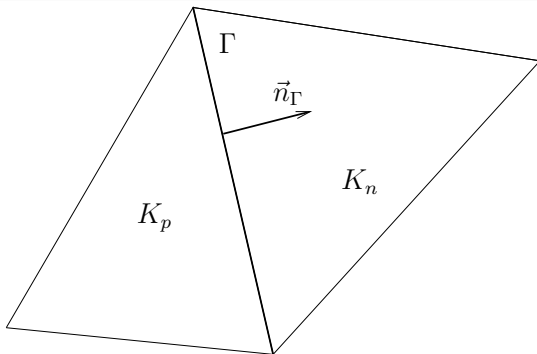
$$S_{hp} \equiv \{v; v \in L^2(\Omega), v|_K \in P_{p_K}(K) \forall K \in \mathcal{T}_h\},$$

- we introduce space  $\mathbf{S}_{hp} = \underbrace{S_{hp} \times \dots \times S_{hp}}_{(d+2) \text{ times}}$

# Example of a function from $S_{hp} \subset H^s(\Omega, \mathcal{T}_h)$



## Notation - trace, mean value, jump



- $v|_\Gamma^{(p)}$   $\equiv$  trace of  $v|_{K_p}$  on  $\Gamma$  and  $v|_\Gamma^{(n)}$   $\equiv$  trace of  $v|_{K_n}$  on  $\Gamma$ ,
- $\langle v \rangle_\Gamma = \frac{1}{2} \left( v|_\Gamma^{(p)} + v|_\Gamma^{(n)} \right)$  and  $[v]_\Gamma = v|_\Gamma^{(p)} - v|_\Gamma^{(n)}$ ,
- $\langle v \rangle_\Gamma \equiv [v]_\Gamma \equiv v|_\Gamma^{(p)}$ ,  $\Gamma \subset \partial\Omega$ .

## DG formulation

- let  $\mathbf{w}$  be a sufficiently regular solution,
- we multiply (1) by  $\varphi \in H^2(\Omega, \mathcal{T}_h)^{d+2}$ ,
- integrate over each  $K \in \mathcal{T}_h$ ,
- apply Green's theorem,
- sum over all  $K \in \mathcal{T}_h$ ,
- we obtain the identity

$$\left( \frac{\partial \mathbf{w}(t)}{\partial t}, \varphi \right) + \tilde{\mathbf{b}}_h(\mathbf{w}(t), \varphi) = 0 \quad \forall \varphi \in H^2(\Omega, \mathcal{T}_h)^{d+2}, \forall t \in (0, T) \quad (2)$$

## Inviscid terms

$$\begin{aligned} \tilde{\mathbf{b}}_h(\mathbf{w}, \varphi) = & \sum_{\Gamma \in \mathcal{F}_h} \int_{\Gamma} \mathbf{H} \left( \mathbf{w}|_{\Gamma}^{(p)}, \mathbf{w}|_{\Gamma}^{(n)}, \vec{n}_{\Gamma} \right) \cdot [\varphi]_{\Gamma} \, dS \\ & - \sum_{K \in \mathcal{T}_h} \int_K \sum_{s=1}^d \mathbf{A}_s(\mathbf{w}) \mathbf{w} \cdot \frac{\partial \varphi}{\partial x_s} \, dx, \end{aligned}$$

where  $\mathbf{H}$  is the Vijayasundaram numerical flux,

$$H \left( \mathbf{w}|_{\Gamma}^{(p)}, \mathbf{w}|_{\Gamma}^{(n)}, \vec{n}_{\Gamma} \right) = \mathbf{P}^+ (\langle \mathbf{w}_h \rangle, \vec{n}) \mathbf{w}_h|_{\Gamma}^{(p)} + \mathbf{P}^- (\langle \mathbf{w}_h \rangle, \vec{n}) \mathbf{w}_h|_{\Gamma}^{(n)}$$

## Space semi-discrete problem

- **method of lines** for the Euler equations,
- $S_{hp} \subset H^2(\Omega, \mathcal{T}_h) \Rightarrow$  identity (2) makes sense for  $\mathbf{w}_h, \varphi_h \in \mathbf{S}_{hp}$ ,
- **approximate solution**  $\mathbf{w}_h(t) \in \mathbf{S}_{hp}$  satisfies the identity:

$$\frac{d}{dt}(\mathbf{w}_h(t), \varphi_h) + \tilde{\mathbf{b}}_h(\mathbf{w}_h(t), \varphi_h) = 0 \quad (3)$$

$$\forall \varphi_h \in \mathbf{S}_{hp}, t \in (0, T),$$

with  $\mathbf{w}_h(0)$  satisfying IC,

- semi-discrete problem (3) represents ODEs,
- **explicit method**  $\Rightarrow$  high restriction on time step,
- full **implicit method**  $\Rightarrow$  system of nonlinear equations

### Semi-implicit method

BDF scheme + linearization of (3)

## Linearization of the inviscid fluxes

- inviscid terms: for  $\tilde{\mathbf{w}}_h, \mathbf{w}_h, \varphi_h \in \mathbf{S}_{hp}$

$$\mathbf{b}_h(\tilde{\mathbf{w}}_h, \mathbf{w}_h, \varphi_h) = - \sum_{K \in \mathcal{T}_h} \int_K \sum_{s=1}^d \mathbf{A}_s(\tilde{\mathbf{w}}_h) \mathbf{w}_h \cdot \frac{\partial \varphi_h}{\partial x_s} \, d\mathbf{x} \\ + \sum_{\Gamma \in \mathcal{F}_h} \int_{\Gamma} \left( \mathbf{P}^+ (\langle \tilde{\mathbf{w}}_h \rangle, \vec{n}) \mathbf{w}_h|_{\Gamma}^{(p)} + \mathbf{P}^- (\langle \tilde{\mathbf{w}}_h \rangle, \vec{n}) \mathbf{w}_h|_{\Gamma}^{(n)} \right) \cdot [\varphi_h] \, dS,$$

- linear with respect to  $\mathbf{w}_h, \varphi_h$ ,
- consistent with  $\tilde{\mathbf{b}}_h$ :

$$\mathbf{b}_h(\mathbf{w}_h, \mathbf{w}_h, \varphi_h) = \tilde{\mathbf{b}}_h(\mathbf{w}_h, \varphi_h) \quad \forall \mathbf{w}_h, \varphi_h \in \mathbf{S}_{hp}.$$



## Higher order semi-implicit BDF-DGM scheme

- partition of  $(0, T) \Rightarrow t_0 < t_1 < \dots < t_r, \tau_k \equiv t_{k+1} - t_k,$

### General higher order scheme

$$\frac{1}{\tau_k} \left( \sum_{l=0}^n \alpha_l \mathbf{w}_h^{k+1-l}, \varphi_h \right) + \mathbf{b}_h \left( \sum_{l=1}^n \beta_l \mathbf{w}_h^{k+1-l}, \mathbf{w}_h^{k+1}, \varphi_h \right) = 0$$

$$\forall \varphi_h \in \mathbf{S}_{hp}, k = n-1, \dots, r-1,$$

- $\alpha_l, \beta_l$  - coefficients of BDF,
- $\mathbf{w}_h^0$  is  $\mathbf{S}_{hp}$ -approximation of  $\mathbf{w}^0$ ,
- $\mathbf{w}_h^l, 1 \leq l \leq n-1$  given by a one-step method.

### Drawback

High order approximation of discontinuous solution  $\Rightarrow$  Gibbs-type oscillations  $\Rightarrow$  instability.

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## Artificial viscosity limiter (1)

**Main idea:** add artificial term to (1) in the form corresponds to the viscous part of the system of the Navier-Stokes equations but with the variable Reynolds number

$$\frac{1}{Re} \Big|_K \approx \mu_{art}(\mathbf{w}_h, K) \Rightarrow \delta_S \sim \mathcal{O}(\mu_{art})$$

- new system  $\sim$  compressible Navier-Stokes equations

$$\frac{\partial \mathbf{w}}{\partial t} + \sum_{s=1}^d \frac{\partial}{\partial x_s} \mathbf{f}_s(\mathbf{w}) = \mu_{art} \sum_{s=1}^d \frac{\partial}{\partial x_s} \mathbf{R}_s(\mathbf{w}, \nabla \mathbf{w}) \quad \text{in } (0, T) \times \Omega$$

- viscous fluxes (without  $Re$ )

$$\mathbf{R}_s(\mathbf{w}, \nabla \mathbf{w}) = \left( 0, \tau_{1s}, \dots, \tau_{ds}, \sum_{r=1}^d \tau_{rs} v_r + \frac{\gamma}{Pr} \frac{\partial \theta}{\partial x_s} \right)^T, \quad s = 1, \dots, d$$

## Artificial viscosity limiter (2)

- viscous part of the stress tensor (without  $Re$ )

$$\tau_{rs} = \left[ \left( \frac{\partial v_s}{\partial x_r} + \frac{\partial v_r}{\partial x_s} \right) - \frac{2}{3} \operatorname{div}(\mathbf{v}) \delta_{rs} \right], \quad r, s = 1, \dots, d,$$

- properties of viscous fluxes  $\mathbf{R}_s$ :

$$\mathbf{R}_s(\mathbf{w}, \nabla \mathbf{w}) = \sum_{k=1}^d \mathbf{K}_{s,k}(\mathbf{w}) \frac{\partial \mathbf{w}}{\partial x_k}, \quad s = 1, \dots, d,$$

where  $\mathbf{K}_{s,k} \in \mathbf{R}^{(d+2) \times (d+2)}$ ,  $s, k = 1, \dots, d$ ,

- add equation of total energy:  $e = c_V \rho \theta + \rho |\mathbf{v}|^2 / 2$ ,

approach  $\equiv$  solution of N.S. equations with "do nothing" BC

## Entropy-based vanishing viscosity (1)

- **nonlinear viscosity**  $\approx$  residual of the entropy equation  $S_{res}(\mathbf{w})$  [Guermond, Pasquetti, 08]
- entropy:  $S = \frac{1}{\gamma-1} \ln \left( \frac{p}{\rho^\gamma} \right)$  ( $\gamma$  - Poisson adiabatic constant)
- entropy form of the energy equation:

$$\frac{\partial \rho S}{\partial t} + \operatorname{div}(\rho S \mathbf{v}) = \frac{D(\mathbf{v})}{\theta} + \mu_{art} \frac{\gamma}{Pr} \frac{\operatorname{div}(\nabla \theta)}{\theta}$$

where  $D(\mathbf{v})$  is dissipation

$$D(\mathbf{v}) = -\frac{2}{3} \mu_{art} (\operatorname{div}(\mathbf{v}))^2 + 2 \mu_{art} \mathbb{D}(\mathbf{v}) \cdot \mathbb{D}(\mathbf{v})$$

with  $\mathbb{D}(\mathbf{v})$  as deformation velocity tensor

## Entropy-based vanishing viscosity (2)

- **discrete entropy residual** (weak formulation):

$$\int_{\Omega} S_{res}(\mathbf{w}) \varphi_h \, dx = \int_{\Omega} \left( \frac{\partial \rho S}{\partial t} + \operatorname{div}(\rho S \mathbf{v}) - \frac{D(\mathbf{v})}{\theta} - \mu_{art} \frac{\gamma}{Pr} \frac{\operatorname{div}(\nabla \theta)}{\theta} \right) \varphi_h \, dx$$

- $\operatorname{supp}(\varphi_h) \subset K$  + Green's theorem imply

$$\begin{aligned} \int_K S_{res}(\mathbf{w})|_K \varphi_h &= \int_K \frac{\partial \rho S}{\partial t} \varphi_h \, dx + \int_{\partial K} \rho S (\mathbf{v} \cdot \vec{n}) \varphi_h \, dS - \int_K \rho S \mathbf{v} \cdot \nabla \varphi_h \, dx \\ &\quad - \int_K \frac{D(\mathbf{v})}{\theta} \varphi_h \, dx - \frac{\gamma}{Pr} \int_{\partial K} \mu_{art}(K) \nabla \theta \cdot \vec{n} \frac{\varphi_h}{\theta} \, dS \\ &\quad + \frac{\gamma}{Pr} \int_K \mu_{art}(K) \nabla \theta \cdot \nabla \left( \frac{\varphi_h}{\theta} \right) \, dx \quad \forall \varphi_h \in P_{pK} \end{aligned}$$

- $S_{res}(\mathbf{w})$  is  $L^2$ -projection onto  $\mathcal{S}_{hp}$ , i.e.  $S_{res}(\mathbf{w})|_K \in P_{pK}$

## Entropy-based vanishing viscosity (3)

- maximal values (to limit viscosity)

$$\mu_{max} = \nu_{max} \max_{K \in \mathcal{T}_h} \rho \quad \text{and} \quad \nu_{max} = \frac{h_K}{\rho_K} \max_{K \in \mathcal{T}_h} \left( |\mathbf{v}| + \sqrt{\gamma \theta} \right)$$

- finally we set

$$\mu_{art}(\mathbf{x})|_K = \min(\mu_{max}, \alpha L h_K \cdot |S_{res}(\mathbf{x})|)$$

with  $h_K$ ...element size,  $L$ ...characteristic length

- regions with **smooth** solution  $\mathbf{w} \Rightarrow S_{res}(\mathbf{w}) \approx 0$
- regions with **shocks**  $\Rightarrow S_{res}(\mathbf{w})$  is large
- simplification: **Laplacian** artificial viscosity

$$\mathbf{K}_{s,k} = \delta_{sk} \mathbf{I} \Rightarrow \sum_{s=1}^d \frac{\partial}{\partial x_s} \mathbf{R}_s(\mathbf{w}, \nabla \mathbf{w}) = \Delta \mathbf{w}$$

- simplification:  $\varphi_h|_K = \chi_K \Rightarrow \mu_{art}$  piecewise constant

## Artificial viscous and penalty terms

$$\begin{aligned} \tilde{\mathbf{a}}_h(\mathbf{w}, \varphi) &= \sum_{K \in \mathcal{T}_h} \int_K \mu_{art}(\mathbf{w}) \left( \sum_{k=1}^d \mathbf{K}_{s,k}(\mathbf{w}) \frac{\partial \mathbf{w}}{\partial x_k} \right) \cdot \nabla \varphi \, dx \\ &- \sum_{\Gamma \in \mathcal{F}_h^I} \int_{\Gamma} \sum_{s=1}^d \left\langle \mu_{art}(\mathbf{w}) \left( \sum_{k=1}^d \mathbf{K}_{s,k}(\mathbf{w}) \frac{\partial \mathbf{w}}{\partial x_k} \right) \right\rangle n_s \cdot [\varphi] \, dS \\ &- \Theta \sum_{\Gamma \in \mathcal{F}_h^I} \int_{\Gamma} \sum_{s=1}^d \left\langle \mu_{art}(\mathbf{w}) \sum_{k=1}^d \mathbf{K}_{k,s}^T(\mathbf{w}) \frac{\partial \varphi}{\partial x_k} \right\rangle n_s \cdot [\mathbf{w}] \, dS, \end{aligned}$$

where  $\Theta = 1$  (SIPG),  $0$  (IIPG),  $-1$  (NIPG)

$$\begin{aligned} \mathbf{J}_h^\sigma(\mathbf{w}, \varphi) &= \sum_{\Gamma \in \mathcal{F}_h^I} \int_{\Gamma} \sigma[\mathbf{w}] \cdot [\varphi] \, dS \quad \text{with } \sigma|_{\Gamma} \equiv \mu_{art} \frac{C_W}{d(\Gamma)}, \quad C_W > 0, \\ d(\Gamma) &= \min(h_{K_p}/p_{K_p}, h_{K_n}/p_{K_n}), \quad \Gamma = K_p \cap K_n \end{aligned}$$



## Linearization of viscous terms

$$\begin{aligned}
 \mathbf{a}_h(\tilde{\mathbf{w}}_h, \mathbf{w}_h, \varphi_h) &= \sum_{K \in \mathcal{T}_h} \int_K \mu_{art} \left( \sum_{k=1}^d \mathbf{K}_{s,k}(\tilde{\mathbf{w}}_h) \frac{\partial \mathbf{w}_h}{\partial x_k} \right) \cdot \nabla \varphi_h \, dx \\
 &\quad - \sum_{\Gamma \in \mathcal{F}'_h} \int_{\Gamma} \sum_{s=1} \left\langle \mu_{art} \left( \sum_{k=1}^d \mathbf{K}_{s,k}(\tilde{\mathbf{w}}_h) \frac{\partial \mathbf{w}_h}{\partial x_k} \right) \right\rangle n_s \cdot [\varphi_h] \, dS \\
 &\quad - \Theta \sum_{\Gamma \in \mathcal{F}'_h} \int_{\Gamma} \sum_{s=1} \left\langle \mu_{art} \sum_{k=1}^d \mathbf{K}_{k,s}^T(\tilde{\mathbf{w}}_h) \frac{\partial \varphi_h}{\partial x_k} \right\rangle n_s \cdot [\mathbf{w}_h] \, dS
 \end{aligned}$$

- **linear** with respect to  $\mathbf{w}_h, \varphi_h$ ,
- **consistency**:

$$\mathbf{a}_h(\mathbf{w}_h, \mathbf{w}_h, \varphi_h) = \tilde{\mathbf{a}}_h(\mathbf{w}_h, \varphi_h) \quad \forall \mathbf{w}_h, \varphi_h \in \mathbf{S}_{hp}$$

## Discrete problem

- we set

$$\mathbf{c}_h(\tilde{\mathbf{w}}_h, \mathbf{w}_h, \varphi_h) := \mathbf{a}_h(\tilde{\mathbf{w}}_h, \mathbf{w}_h, \varphi_h) + \mathbf{b}_h(\tilde{\mathbf{w}}_h, \mathbf{w}_h, \varphi_h) + \mathbf{J}_h(\mathbf{w}_h, \varphi_h),$$

- partition  $(0, T) \rightarrow t_0 < t_1 < \dots < t_r$ ,  $\tau_k \equiv t_{k+1} - t_k$ ,

### Shock-capturing scheme

$$\frac{1}{\tau_k} \left( \sum_{l=0}^n \alpha_l \mathbf{w}_h^{k+1-l}, \varphi_h \right) + \mathbf{c}_h \left( \sum_{l=1}^n \beta_l \mathbf{w}_h^{k+1-l}, \mathbf{w}_h^{k+1}, \varphi_h \right) = 0$$

$$\forall \varphi_h \in \mathbf{S}_{hp}, \quad k = n-1, \dots, r-1,$$

- $\alpha_l, \beta_l$  - coefficients of BDF,
- $\mathbf{w}_h^0$  is  $\mathbf{S}_{hp}$ -approximation of  $\mathbf{w}^0$ ,
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## Basis $\mathbf{S}_{hp}$

- local basis:

$$B_i = \{\psi_{ij}, \psi_{ij} \in \mathbf{S}_{hp}, \text{supp}(\psi_{ij}) \subset K_i, j = 1, \dots, \text{dof}_i\}, \\ i = 1, \dots, \#\mathcal{T}_h,$$

- global basis:

$$B = \{\psi_{ij}, \psi_{ij} \in B_i, j = 1, \dots, \text{dof}_i, i = 1, \dots, \#\mathcal{T}_h\},$$

### Linear algebraic representation

$$\mathbf{w}_h^k(x) = \sum_{K_i \in \mathcal{T}_h} \sum_{j=1}^{\text{dof}_i} \xi_{kij} \psi_{ij}, \quad x \in \Omega, \quad k = 0, 1, \dots, r,$$

$$\mathbf{w}_h^k(x) \leftrightarrow \mathbf{W}_k = \{\xi_{kij}\}_{ij} \in \mathbf{R}^{DOF}, \quad DOF = \sum_{K_i \in \mathcal{T}_h} \text{dof}_i.$$

## Linear algebraic problem

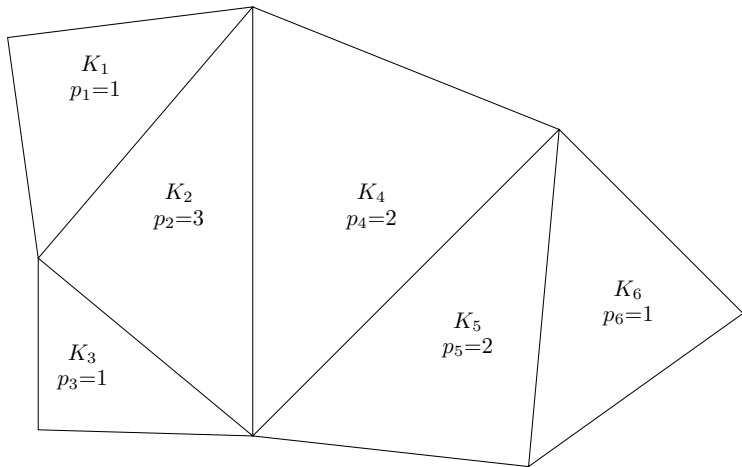
### BDF-DG scheme

$$(\mathbf{M} + \tau_k \mathbf{C}_k(\mathbf{W}_k)) = \mathbf{q}_k, \quad k = 1, \dots, r,$$

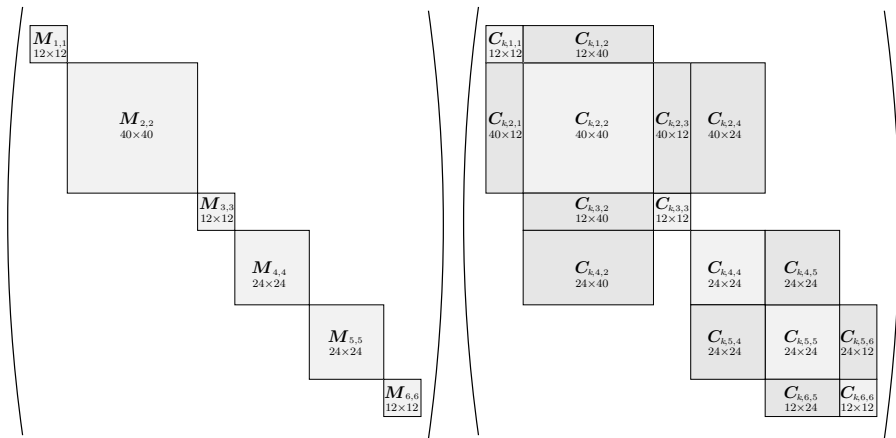
where

- $\mathbf{W}_k$  unknown vector,
- $\mathbf{M}$  mass matrix,
- $\mathbf{C}_k$  "flux" matrix representing inviscid, viscous and penalty terms,
- $\mathbf{q}_k$  right-hand side (BC),
- $\tau_k$  time step.

## Fictional triangulation



# Matrix structure of $\mathbf{M}$ and $\mathbf{C}_k$



## Other implementation aspects

- restarted GMRES with block diagonal preconditioning,
- orthonormal basis of  $S_{hp} \Rightarrow \mathbf{M} =$  identity matrix,
- time step for semi-implicit scheme: adaptive and/or heuristic choice,
- penalty parameter  $\sigma$  is set according to NIPG variant:  
 $C_W = 1.0$ ,
- 2D Fortran code: program ADGFEM [Dolejší]  
(KNM MFF UK)

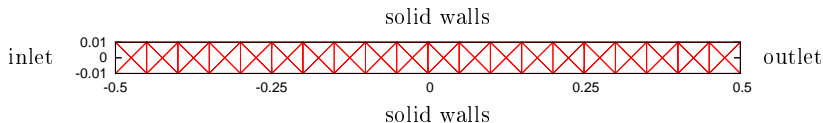


# Lax problem (1)

- 1D test case: shock, contact discontinuity, rarefaction wave
- computational domain  $[-0.5, 0.5]$ ,
- initial conditions:

$$(\rho, v_1, p) = \begin{cases} (0.445, & 0.698, & 3.528) & \text{if } x \leq 0, \\ (0.5, & 0.0, & 0.571) & \text{if } x > 0. \end{cases}$$

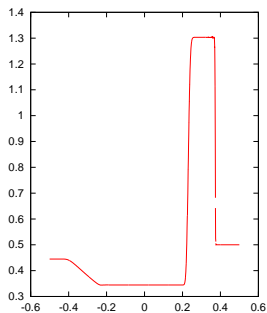
- reflecting boundary conditions,
- interval  $[-0.5, 0.5]$  replaced with 2D computational domain



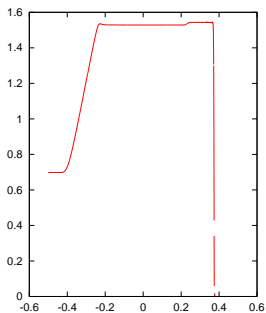
- added BC on solid walls:  $\mathbf{v} \cdot \vec{n} = 0 \Leftrightarrow v_2 = 0$

## Lax problem (2)

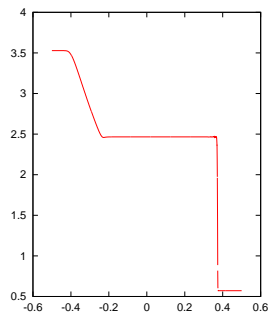
- $P_3$  approximation, BDF 1<sup>st</sup> order
- uniform grid with 2 000 elements, final time  $T = 0.15$ ,
- distribution of density, velocity and pressure



density



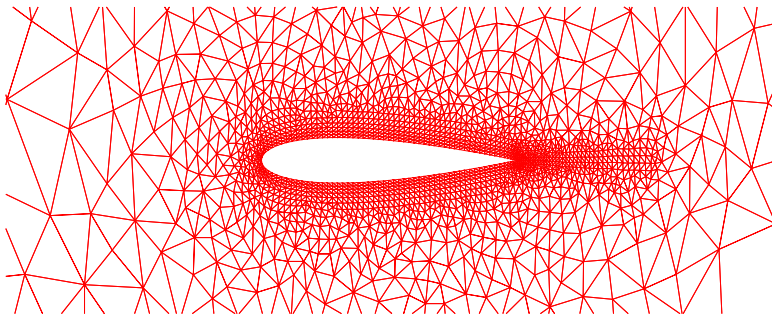
velocity



pressure

## NACA 0012 profile - steady flow (1)

- **steady transonic inviscid** flow around the NACA0012 airfoil
- $M = 0.8$ ,  $\alpha = 1.25^\circ$  (non-symmetric flow),
- triangular adaptive refined grid having 4 544 elements,
- adaptive BDF scheme,  $P_1 - P_3$  approximation



## NACA 0012 profile - steady flow (2)

- test case from project **ADIGMA**:  
A EUROPEAN PROJECT ON THE DEVELOPMENT OF  
**AD**APTIVE **HIG**HER ORDER VARIATIONAL **M**ETHODS  
FOR **A**EROSPACE APPLICATIONS
- coefficients of drag  $c_D$  and lift  $c_L$  (comparison with UNST)

method	$c_D$	$c_L$	$DOF$
BDF-DGM – $P_1$	0.02426	0.33684	54 528
BDF-DGM – $P_2$	0.02300	0.34065	109 056
BDF-DGM – $P_3$	0.02277	0.35587	181 760
NSDG2D – $P_3$ [UNST]	0.02254	0.36448	181 760
NSDG2D – $P_5$ [UNST]	0.02276	0.35366	381 696

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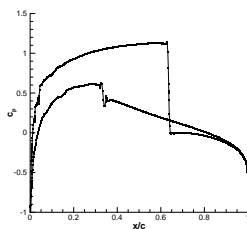
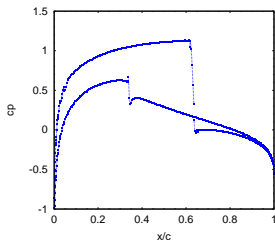
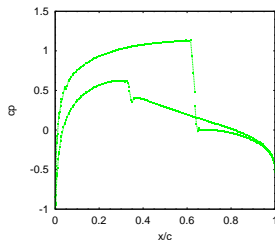
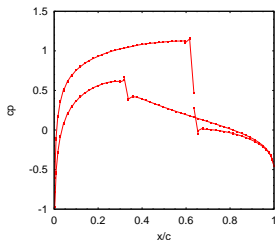
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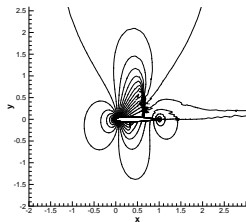
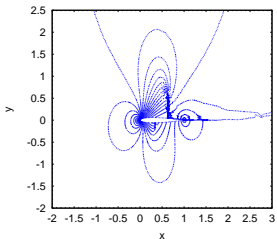
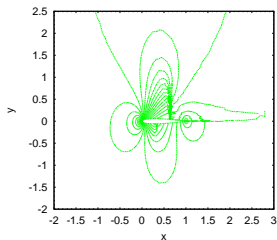
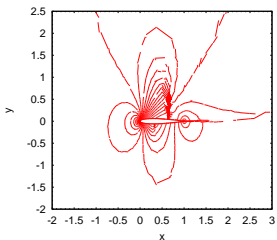
# NACA 0012 profile - pressure coefficient

- comparison of  $P_1$ ,  $P_2$ ,  $P_3$  and  $P_5$  (UNST) approximations



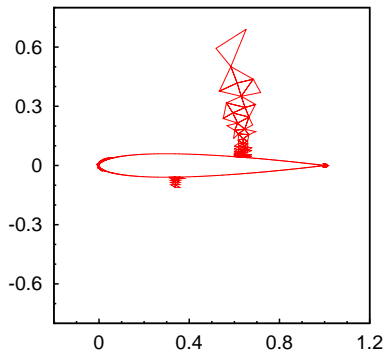
# NACA 0012 profile - Mach number isolines

- comparison of  $P_1$ ,  $P_2$ ,  $P_3$  and  $P_5$  (UNST) approximations

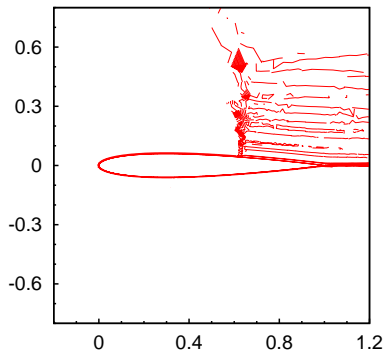




# NACA 0012 profile - artificial viscosity vs. entropy

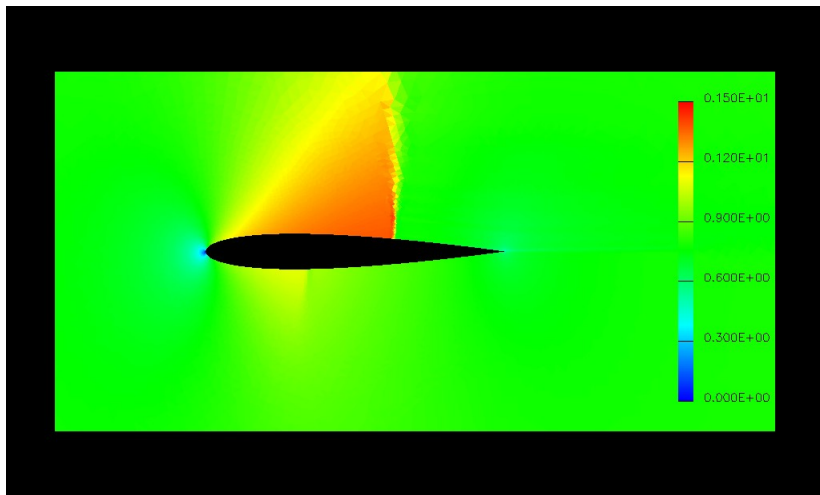


$\mu_{art} > 10^{-3}$

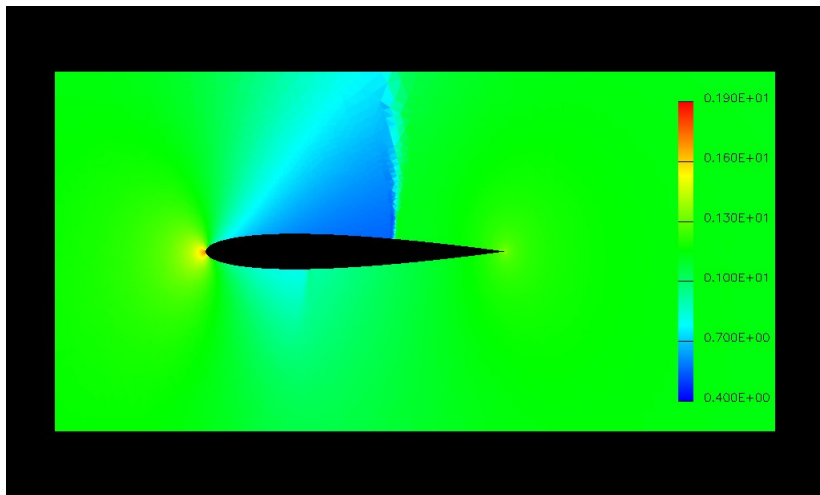


entropy

# NACA 0012 - Mach number distribution, $t \rightarrow \infty$

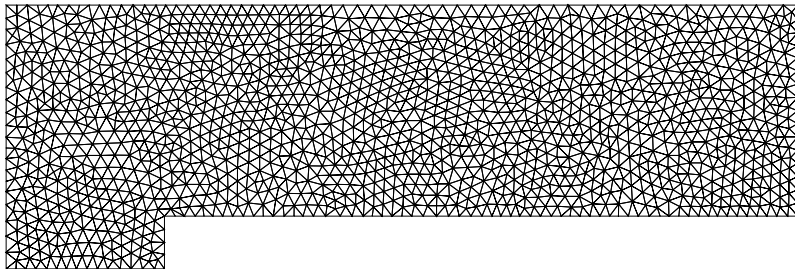


# NACA 0012 - pressure distribution, $t \rightarrow \infty$



## Forward facing step - unsteady flow

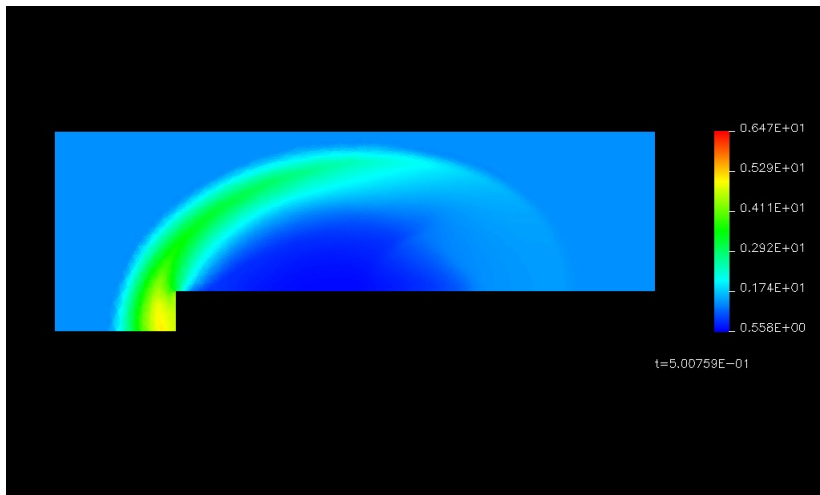
- benchmark of **unsteady inviscid** flow
- IC:  $\mathbf{w} = (1.4, 3, 0, 1)$ , inflow/outflow BC form IC,
- simulation for  $t \in (0, 4)$ ,
- grid with 2914 triangles,  $P_2$  approximation, **2<sup>nd</sup> order** in time



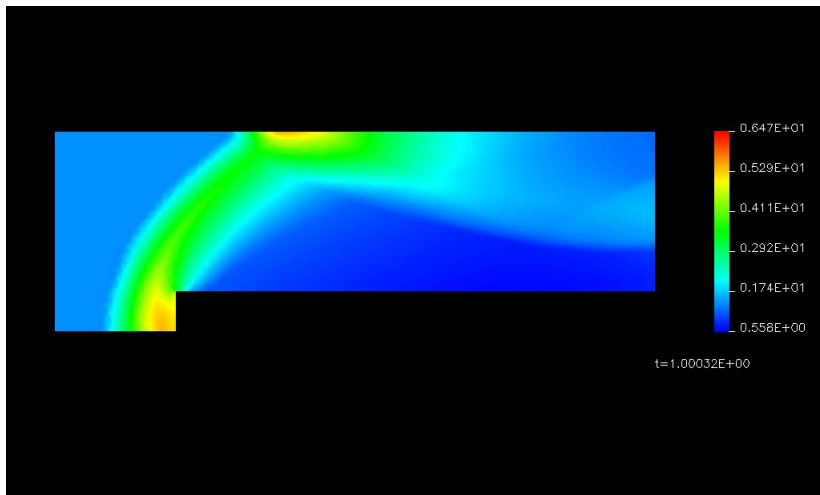
## FFS - density distribution, $t=0.0$



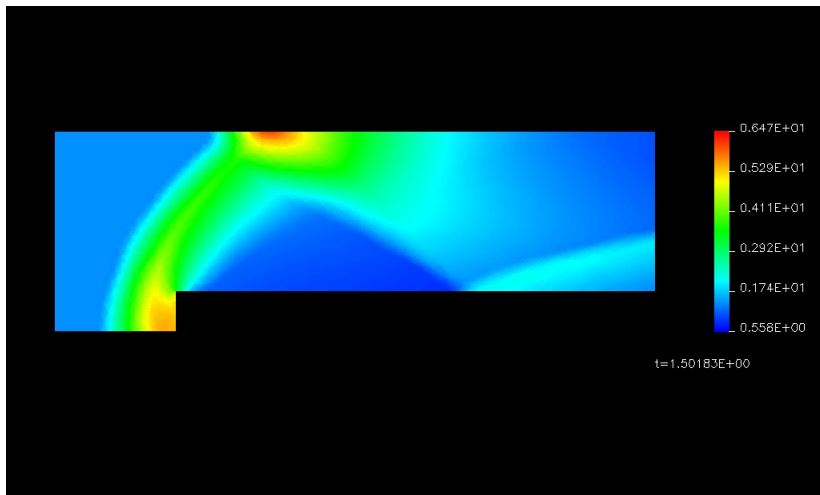
## FFS - density distribution, $t=0.5$



## FFS - density distribution, $t=1.0$

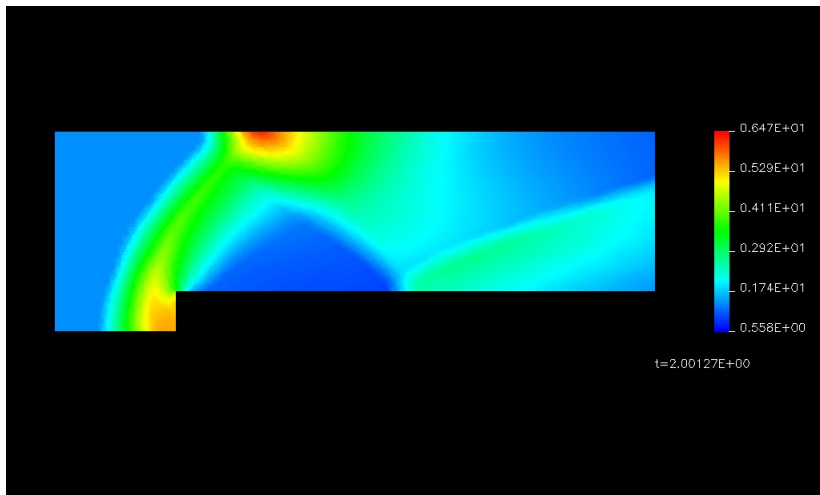


## FFS - density distribution, $t=1.5$

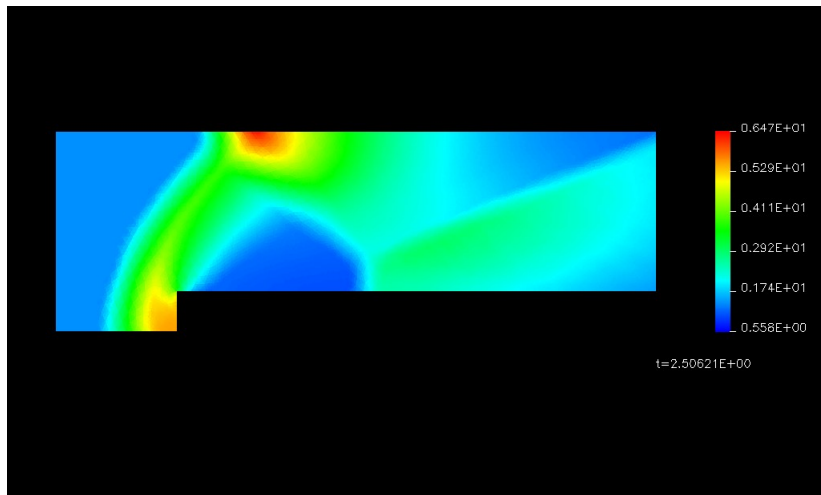




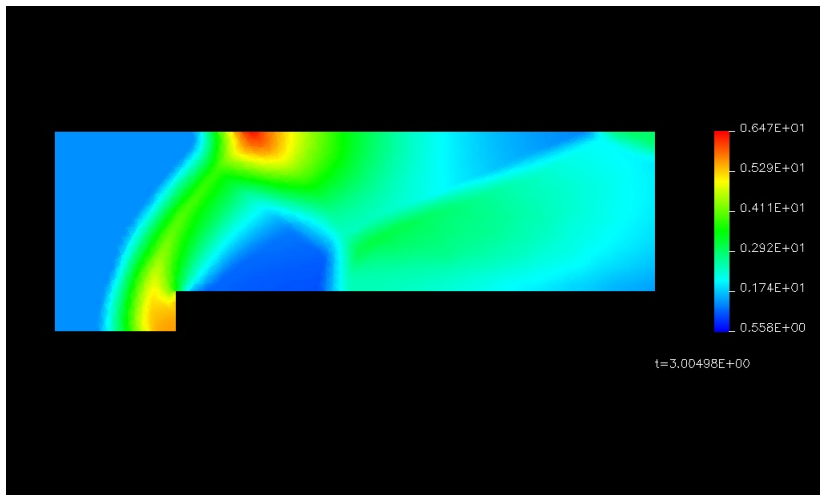
## FFS - density distribution, $t=2.0$



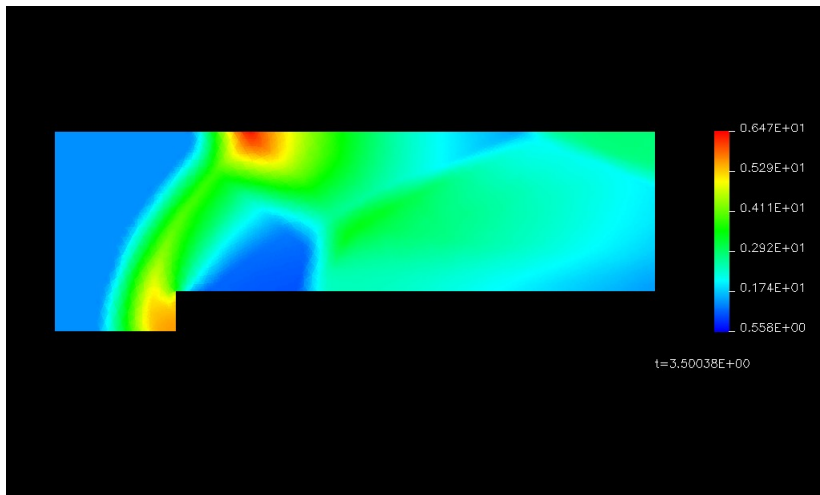
## FFS - density distribution, $t=2.5$



## FFS - density distribution, $t=3.0$



## FFS - density distribution, $t=3.5$



## FFS - density distribution, $t=4.0$

