

# Numerical approach to a rate-independent model of decohesion in laminated composites

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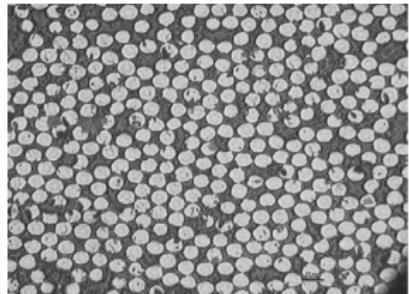
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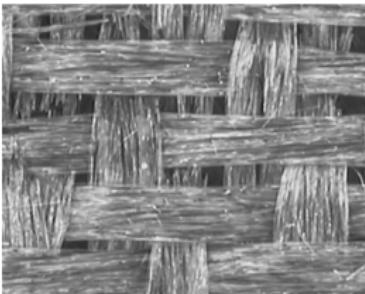
PANM 15: Programy a algoritmy numerické matematiky 15  
6.–11. června, Dolní Maxov

# Interfaces in composite materials

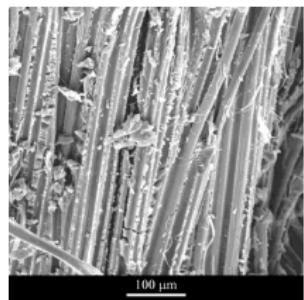
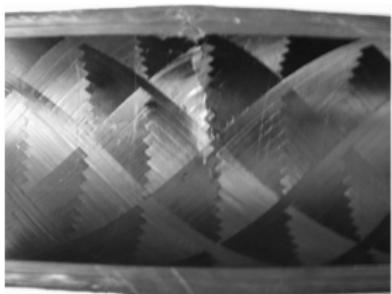
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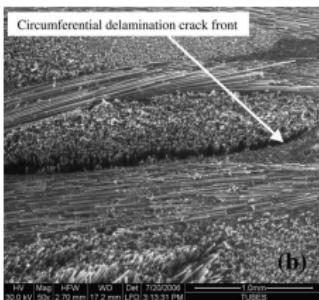
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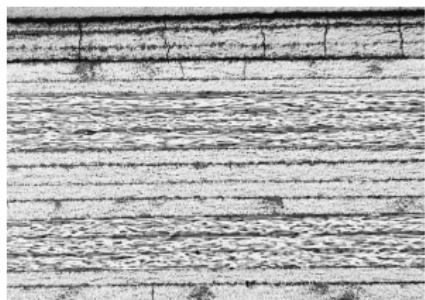
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$\approx 10 \mu\text{m}$



$\approx 1 \text{ mm}$



$\approx 1 \text{ dm}$

Essential aspect for reliable design of composite structures

# Modelling assumptions

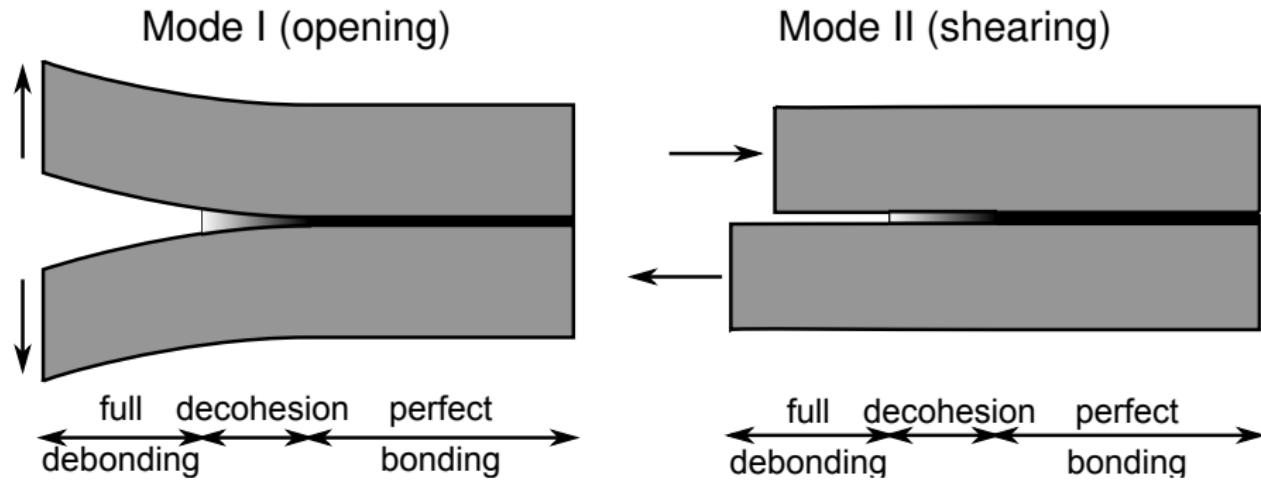
(Macroscopic delamination)



- Energy-based framework
- Decohesion  $\equiv$  displacement jumps at interfaces
- Inelastic phenomena concentrated at interfaces
- Rate-independent (quasi-static) approximation
- Frictionless contact conditions
- Small-strain setting

# Basic concepts of interfacial damage mechanics

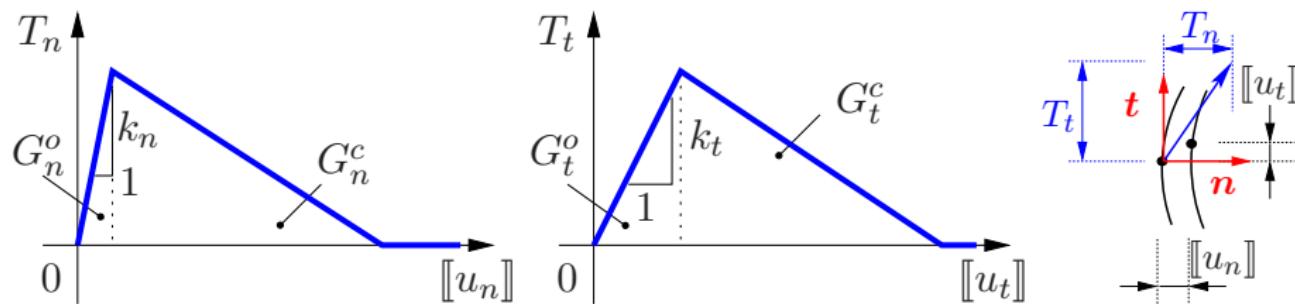
## Constitutive characterization



- Initial stiffnesses:  $k_n, k_t$  ( $\text{Nm}^{-3}$ )
- Activation energies:  $G_n^o, G_t^o$  ( $\text{Jm}^{-2}$ )
- Dissipated energies:  $G_n^c, G_t^c$  ( $\text{Jm}^{-2}$ )
- Mode interaction parameters

# Basic concepts of interfacial damage mechanics

## State variables



- Displacement jump decomposition

$$\llbracket u_n \rrbracket = \llbracket \mathbf{u} \rrbracket \cdot \mathbf{n} \geq 0 \quad \llbracket u_t \rrbracket = \| \llbracket \mathbf{u} \rrbracket - \llbracket u_n \rrbracket \mathbf{n} \| \quad (1)$$

- Intensity of adhesion

$$T_n = \zeta k_n \llbracket u_n \rrbracket \quad T_t = \zeta k_t \llbracket u_t \rrbracket \quad 0 \leq \zeta \leq 1$$

# Basic concepts of interfacial damage mechanics

Interaction criteria (VALOROSO & CHAMPANEY, 2006)

- Mode mixity parameter

$$\psi(\llbracket \mathbf{u} \rrbracket) = \sqrt{\frac{k_t}{k_n}} \frac{\llbracket u_t \rrbracket}{\llbracket u_n \rrbracket} \quad (\text{mode I} =) \quad 0 \leq \psi < +\infty \quad (= \text{mode II})$$

- Mode interaction criteria

$$\left( \frac{G^o(\psi)}{(1 + \psi^2)G_n^o} \right)^{a_1} + \left( \frac{\psi^2 G^o(\psi)}{(1 + \psi^2)G_t^o} \right)^{a_2} = 1$$

$$\left( \frac{G^c(\psi)}{(1 + \psi^2)G_n^c} \right)^{b_1} + \left( \frac{\psi^2 G^c(\psi)}{(1 + \psi^2)G_t^c} \right)^{b_2} = 1$$

parameters  $a_1, a_2, b_1, b_2$  fitted from experiments or taken as 2

# Motivation

Challenges in modelling delamination phenomena

## Theoretical aspects

- ① No supporting existence theory for mixed-mode description
- ② (At least) abstract approximation results

## Computational aspects

- ① Oscillations of interfacial tractions for almost perfect interfaces  
 $(k_\bullet \rightarrow \infty)$
- ② Unstable response for brittle interfaces, leading to oscillation of quantity of interests for coarse meshes
- ③ Efficient and reliable resolution of frictionless contact

# Outline

- 1 Energetic rate-independent systems
  - Energetic solution
  - Incremental energy minimization
  - Existence and approximation results
  - Application to delamination problem
- 2 Numerical treatment
  - Incremental optimization problems
  - Alternate minimization
  - Backtracking strategy
  - Application of FETI-based methods
- 3 Examples
- 4 Summary and outlook

# Energetic rate-independent systems

Notation (MIELKE, LEVITAS & THEIL, 2002; MIELKE & ROSSI, 2007)

- Domain  $\Omega \subset \mathbb{R}^d$ , time interval  $\mathbb{I} = [0; T]$
- State variables:  $\mathbf{q} = (\mathbf{u}, \mathbf{z}) \in \mathcal{Q} = \mathcal{U} \times \mathcal{Z}$ 
  - Displacements  $\mathbf{u} : \Omega \rightarrow \mathbb{R}^d$
  - Internal variables  $\mathbf{z} : \Omega \rightarrow \mathbb{R}^m$
- Constitutive description
  - Stored energy  $E(t, \mathbf{q}) : \mathbb{I} \times \mathcal{Q} \rightarrow \overline{\mathbb{R}}$
  - Dissipation rate  $D(\mathbf{q}, \dot{\mathbf{z}}) : \mathcal{Q} \times \mathcal{Q} \rightarrow [0; +\infty]$   
**State-dependent**  
positively 1-homogeneous  $\rightarrow$  rate-independence
- Stable sets

$$\mathcal{S}(t) = \{\mathbf{q} \in \mathcal{Q} : E(t, \mathbf{q}) \leq E(t, \tilde{\mathbf{q}}) + D(\mathbf{q}, \tilde{\mathbf{z}} - \mathbf{z}) \text{ for all } \tilde{\mathbf{q}} \in \mathcal{Q}\}$$

# Energetic rate-independent systems

Energetic solution (MIELKE and co-workers, 2002—)

## Energetic solution

Given  $\mathbf{q}(0) \in \mathcal{Q}$ , the process  $\mathbf{q} : \mathbb{I} \rightarrow \mathcal{Q}$  is an *energetic solution* to the rate-independent system  $(\mathcal{Q}, E, D)$  if for all  $t \in \mathbb{I}$  we have

- **Stability**

$$\mathbf{q}(t) \in \mathcal{S}(t)$$

- **Energy inequality**

$$E(t, \mathbf{q}(t)) + \int_0^t D(\mathbf{q}(s), \dot{\mathbf{z}}(s)) \, ds \leq E(0, \mathbf{q}(0)) + \int_0^t \partial_s E(s, \mathbf{q}(s)) \, ds$$

- Available existence results either assume  $D(\mathbf{q}, \dot{\mathbf{z}}) = D(z, \dot{z})$  or rely on a uniform convexity of  $E(t, \cdot)$  (MR, 2007)

# Energetic rate-independent systems

Incremental energy minimization (MIELKE and co-workers, 2002—)

- Rothe method

$$\mathcal{P}_\tau = \{0 = t_\tau^0 < t_\tau^1 < \dots < t_\tau^M = T\} \quad \text{with} \quad \tau = \max_{j=1,\dots,M} \{t_\tau^j - t_\tau^{j-1}\}$$

## Incremental problems

Given  $\mathbf{q}_\tau^0 = \mathbf{q}(0) \in \mathcal{Q}$ , solve for  $k = 1, \dots, M$

$$\mathbf{q}_\tau^k \in \arg \min_{\mathbf{q} \in \mathcal{Q}} E(t_k, \mathbf{q}) + D(\mathbf{q}_\tau^{k-1}, \mathbf{z} - \mathbf{z}_\tau^{k-1}) \quad (2)$$

- Used to show convergence as  $\tau \rightarrow 0$
- Provides convenient basis for numerical treatment (ORTIZ & STAINIER, 1999 —)

# Energetic rate-independent systems

Existence results (MIELKE and co-workers, 2002—; KRUŽÍK & ZEMAN, in preparation)

## Assumptions (Stored energy)

- $E(t, \cdot)$  is weakly lower semicontinuous and coercive

$$E(t, \mathbf{q}) < \infty \Rightarrow |\partial_t E(t, \mathbf{q})| \leq C_0(C_1 + E(t, \mathbf{q}))$$

- $\partial_t E(t, \cdot)$  is weakly continuous for all  $t \in \mathbb{I}$  and satisfies

$$|\partial_t E(t_1, \mathbf{q}) - \partial_t E(t_2, \mathbf{q})| \leq \omega(|t_1 - t_2|)$$

for a non-decreasing  $\omega : \mathbb{I} \rightarrow \mathbb{R}_0^+$

# Energetic rate-independent systems

Existence results (MIELKE and co-workers, 2002—; KRUŽÍK & ZEMAN, in preparation)

## Assumptions (Dissipation + “Compatibility”)

- For all  $z \in \mathcal{Z}$  and  $q_1, q_2, q_3 \in \mathcal{Q}$

$$C_0 \|z_1 - z_2\|_{\mathcal{X}} \leq D(q_1, z_2 - z_1)$$

$$|D(q_1, z) - D(q_2, z)| \leq C_1 \|q_2 - q_1\|_{\mathcal{Q}} \|z\|_{\mathcal{Z}}$$

$$D(q_1, z_3 - z_1) \leq D(q_1, z_2 - z_1) + D(q_1, z_3 - z_2)$$

where  $\mathcal{Z} \Subset \mathcal{X}$

- For each sequence  $\{(t_n, q_n)\}_{n \in \mathbb{N}}$  with  $q_n \in \mathcal{S}(t_n)$  and  $t_n \rightarrow t^*$  and  $q_n \rightharpoonup q^*$  and  $\sup_n E(t_n, q_n) < \infty$  it holds

$$q_* \in \mathcal{S}(t^*) \quad \partial_t E(t^*, q^n) \rightarrow \partial_t E(t^*, q^*)$$

# Energetic rate-independent systems

Existence results (MIELKE and co-workers, 2002—; KRUŽÍK & ZEMAN, in preparation)

- Solution of the incremental problem exists for all  $k$  and  $\tau$  and is stable
- It satisfies a-posteriori two-sided energy estimates

$$\begin{aligned} \int_{t_\tau^{k-1}}^{t_\tau^k} \partial_s E(s, \mathbf{q}_\tau^k) \, ds &\leq E(t_\tau^k, \mathbf{q}_\tau^k) - E(t_\tau^{k-1}, \mathbf{q}_\tau^{k-1}) \\ &+ D(\mathbf{q}_\tau^{k-1}, \mathbf{z}_\tau^k - \mathbf{z}_\tau^{k-1}) \leq \int_{t_\tau^{k-1}}^{t_\tau^k} \partial_s E(s, \mathbf{q}_\tau^{k-1}) \, ds \end{aligned} \quad (3)$$

- As  $\tau \rightarrow 0$ , there exist  $\mathbf{z} \in BV(\mathbb{I}; \mathcal{X})$  and  $\mathbf{u}(t) : \mathbb{I} \rightarrow \mathcal{U}$  such that

$$(\mathbf{u}(t), \mathbf{z}(t)) \in \mathcal{S}(t) \quad \text{and} \quad \hat{\mathbf{z}}_\tau \rightarrow \mathbf{z} \text{ in } BV(\mathbb{I}; \mathcal{X})$$

and the process satisfies the energy inequality.

# Energetic rate-independent systems

Abstract approximation result (MIELKE & ROUBÍČEK, 2009)

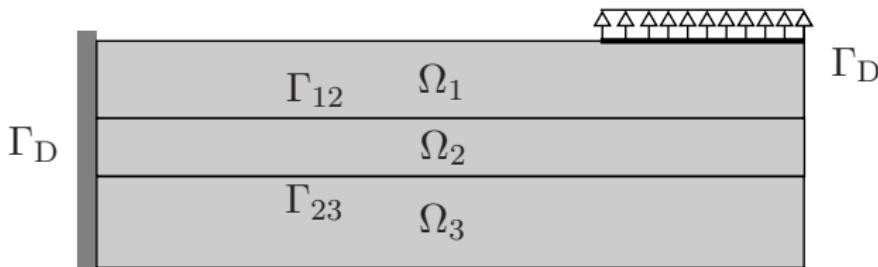
## Energetic density of $\mathcal{Q}_h \subset \mathcal{Q}$

For each  $(t, \mathbf{q}) \in \mathbb{I} \times \mathcal{Q}$ , there exists  $\{\mathbf{q}_h\}$  such that  $\mathbf{q}_h \in \mathcal{Q}_h$ ,  $\mathbf{q}_h \rightarrow \mathbf{q}$  and  $E(t, \mathbf{q}_h) \rightarrow E(t, \mathbf{q})$

- Incremental energy minimization yields stable approximate solution  $\{\mathbf{q}_{\tau,h}^k\}$
- Approximate solutions  $\{\mathbf{q}_{\tau,h}^k\}$  satisfy the two-sided energy estimate
- If  $\mathcal{Q}_h$  satisfies adjusted compatibility condition, the solution converges to an energetic solution as  $\tau \rightarrow 0$  and  $h \rightarrow 0$

# Application to delamination problem

Geometry (KOČVARA, MIELKE & ROUBÍČEK, 2006)



- Body  $\Omega \subset \mathbb{R}^d$  with the Lipschitz boundary  $\Gamma$
- Time-dependent Dirichlet boundary conditions at  $\Gamma_D$
- Collection of disjoint bodies  $\Omega_\alpha$  with Lipschitz boundaries  $\Gamma_\alpha$ ,  
 $\alpha = 1, 2, \dots, m$  ( $\Gamma_{D,\alpha} = \Gamma_\alpha \cap \Gamma_D \neq \emptyset$ )
- Internal boundaries  $\Gamma_{\alpha\beta} = \Gamma_\alpha \cap \Gamma_\beta$  ( $\alpha > \beta$ ) with normal vectors  
 $n_{\alpha\beta}$

# Application to delamination problem

State variables and data (KOČVARA, MIELKE & ROUBÍČEK, 2006; KRUŽÍK & ZEMAN)

- Domain displacements  $\mathbf{u}_\alpha : \Omega_\alpha \rightarrow \mathbb{R}^d$
- Internal variable  $\zeta_{\alpha\beta} : \Gamma_{\alpha\beta} \rightarrow \mathbb{R}$
- Function spaces

$$\begin{aligned}\mathcal{U} &= \prod_{\alpha=1}^m \left\{ \mathbf{u}_\alpha \in W^{1,2}(\Omega_\alpha; \mathbb{R}^d), \mathbf{u}_\alpha|_{\Gamma_{D\alpha}} = \mathbf{0} \right\} \\ \mathcal{Z} &= \prod_{\alpha=1}^m \prod_{\beta=\alpha+1}^m \left\{ \zeta_{\alpha\beta} \in W^{1,2}(\Gamma_{\alpha\beta}) \right\} \\ \mathcal{X} &= \prod_{\alpha=1}^m \prod_{\beta=\alpha+1}^m \left\{ \zeta_{\alpha\beta} \in L^1(\Gamma_{\alpha\beta}) \right\}\end{aligned}$$

- Hard-device loading  $\mathbf{w}_D \in C^1(\mathbb{I}; W^{1/2,2}(\Gamma_D; \mathbb{R}^d))$

# Application to delamination problem

Stored energy (KOČVARA, MIELKE & ROUBÍČEK, 2006; KRUŽÍK & ZEMAN)

- Stored energy

$$E(t, \mathbf{q}) = \sum_{\alpha=1}^m E_\alpha(t, \varepsilon(\mathbf{u}_\alpha)) + \sum_{\alpha=1}^m \sum_{\beta=\alpha+1}^m E_{\alpha\beta}(\mathbf{u}_\alpha - \mathbf{u}_\beta, \zeta_{\alpha\beta})$$

- Bulk energy

$$E_\alpha(t, \varepsilon) = \frac{1}{2} \int_{\Omega_\alpha} \mathbf{C}_\alpha (\varepsilon + \varepsilon(\mathbf{u}_{D,\alpha}(t))) : (\varepsilon + \varepsilon(\mathbf{u}_{D,\alpha}(t))) \, d\Omega$$

- Interfacial energy from (1)

$$E_{\alpha\beta}([\![\mathbf{u}]\!], \zeta) = \begin{cases} \int_{\Gamma_{\alpha\beta}} \frac{\zeta}{2} \mathbf{k}_{\alpha\beta} [\![\mathbf{u}]\!] \cdot [\![\mathbf{u}]\!] + f_{\alpha\beta}(\psi([\![\mathbf{u}]\!]), \zeta) + \frac{\kappa}{2} (\zeta')^2 \, d\Gamma \\ \text{for } [\![\mathbf{u}]\!]_n \geq 0 \text{ and } 0 \leq \zeta \leq 1 \text{ a.e. on } \Gamma_{\alpha\beta} \\ +\infty \text{ otherwise} \end{cases}$$

# Application to delamination problem

Dissipation function (KOČVARA, MIELKE & ROUBÍČEK, 2006; KRUŽÍK & ZEMAN)

- Dissipation rate

$$D(\mathbf{q}, \dot{\mathbf{z}}) = \sum_{\alpha=1}^m \sum_{\beta=\alpha+1}^m D_{\alpha\beta}(\mathbf{u}_\alpha - \mathbf{u}_\beta, \zeta_{\alpha\beta}, \dot{\zeta}_{\alpha\beta})$$

- Interfacial dissipation

$$D_{\alpha\beta}([\![\mathbf{u}]\!], \zeta, \dot{\zeta}) = \begin{cases} - \int_{\Gamma_{\alpha\beta}} G_{\alpha\beta}^c(\psi([\![\mathbf{u}]\!])) \dot{\zeta} \, d\Gamma \text{ for } \dot{\zeta} \leq 0 \\ +\infty \text{ otherwise} \end{cases}$$

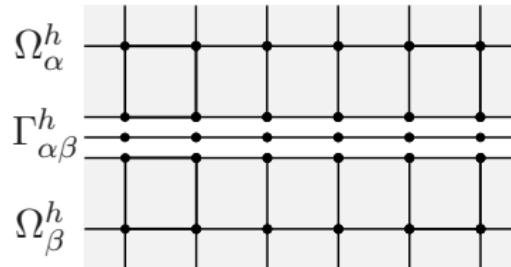
- (KMR, 2006)  $\Rightarrow$  Existence of energetic solution
- (MR, 2009)  $\Rightarrow$  Approximation results for state-independent dissipation and for discretization with  $\mathcal{U}$  with  $P_1$  elements and  $\mathcal{Z}$  with  $P_0$  elements

# Numerical treatment

## Strategy and algebraic re-formulation

- General strategy

- Reduce numerics to interfaces → FETI-family methods  
(DOSTÁL, 2009)



- Discrete state variables

- Nodal displacements:  $u_\alpha$
- Interfacial displacement jumps:  $\llbracket u_{\alpha\beta} \rrbracket$
- Adhesion intensity variables:  $z_{\alpha\beta}$

$$(u, \llbracket u \rrbracket) \in \mathbb{K}^u = \{(v, \llbracket v \rrbracket) : B_d v = 0, B_e v = \llbracket v \rrbracket, B_i \llbracket v \rrbracket \geq 0\}$$

$$z \in \mathbb{K}^z = \{y : 0 \leq y \leq 1\}$$

# Numerical treatment

## Discrete energy functions

- Energy stored in domains

$$E_{\Omega,h}(t, \mathbf{u}) = \frac{1}{2} (\mathbf{u} + \mathbf{u}_D(t))^T \mathbf{K} (\mathbf{u} + \mathbf{u}_D(t)) \quad \mathbf{K} = \begin{bmatrix} \mathbf{K}_1 & & \\ & \ddots & \\ & & \mathbf{K}_m \end{bmatrix}$$

- Energy stored at interfaces

$$E_{\Gamma,h}([\![\mathbf{u}]\!], z) = \frac{1}{2} [\![\mathbf{u}]\!]^T \mathbf{k}(z) [\![\mathbf{u}]\!] + f_h(\psi([\![\mathbf{u}]\!]), z)$$

$f_h(\psi, \cdot)$  convex in  $z$

- Interfacial dissipation:  $z \leq 0$

$$D_h([\![\mathbf{u}]\!], z) = -g(\psi([\![\mathbf{u}]\!]))^T z$$

# Numerical treatment

## Incremental optimization problems

- Incremental energy (2)

$$\begin{aligned} I_{k,h}(z, u, [\![u]\!]) &= \frac{1}{2} u^T K u + u^T K u_D(t_k) + \frac{1}{2} [\![u]\!]^T k(z) [\![u]\!] + f_h(\psi([\![u]\!]), z) \\ &\quad - z^T g(\psi([\![u]\!]_{k-1})) \end{aligned}$$

## Incremental optimization problems

Given  $(z_0, u_0, [\![u]\!]_0) \in \mathbb{K}^z \times \mathbb{K}^u$  solve for  $k = 1, \dots, M$

$$(z_k, u_k, [\![u]\!]_k) = \arg \min_{0 \leq z \leq z_{k-1}} \min_{(u, [\![u]\!]) \in \mathbb{K}^u} I_{k,h}(z, u, [\![u]\!])$$

- Large-scale, sparse, separately convex and constrained problem

# Numerical treatment

Incremental minimization (BOURDIN, FRANCFORST & MARIGO, 2000; BOURDIN, 2007)

- Alternate minimization algorithm ( $\equiv$  staggered scheme)

## Alternate minimization algorithm

1: Set  $j = 0$ ,  $z^{(0)} = z_{k-1}$ ,  $u^{(0)} = u_{k-1}$ ,  $\llbracket u \rrbracket^{(0)} = \llbracket u \rrbracket_{k-1}$

2: **repeat**

3:     Set  $j = j + 1$

4:     Solve  $(u^{(j)}, \llbracket u \rrbracket^{(j)}) = \arg \min_{(u, \llbracket u \rrbracket) \in \mathbb{K}^u} I_{k,h}(z^{(j-1)}, u, \llbracket u \rrbracket)$

5:     Solve  $z^{(j)} = \arg \min_{0 \leq z \leq z_{k-1}} I_{k,h}(z, u^{(j)}, \llbracket u \rrbracket^{(j)})$

6: **until**  $\|z^{(j)} - z^{(j-1)}\|_\infty \leq \delta$

7: Set  $u_k = u^{(j)}$ ,  $\llbracket u \rrbracket_k = \llbracket u \rrbracket^{(j)}$ ,  $z_k = z^{(j)}$

- Step 5 can be performed element-by-element at interface
- Step 4 well-suited for FETI-based solvers

# Numerical treatment

Backtracking strategy (MIELKE, ROUBÍČEK & ZEMAN, 2010; BENEŠOVÁ, 2009–)

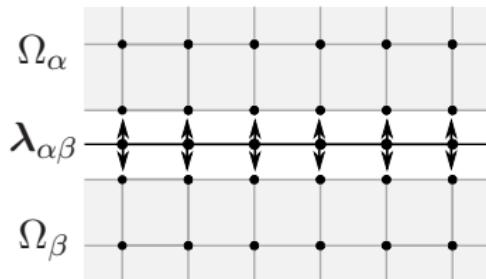
- Convergence to a *critical point* of the objective function
- Not good enough – theory relies on the *global* minimization

## Backtracking algorithm

```
1 : Set  $k = 1$ ,  $z_{-1} = 1$ ,  $z_0 = 1$ ,  $z^{(0)} = 1$ 
2 : repeat
3 :   Determine  $z_k$  using the alternate minimization algorithm
   :   for time  $t_k$  and initial value  $z^{(0)}$ .
4 :   Set  $z^{(0)} = z_k$ 
5 :   if two-sided estimate (3) is satisfied with tolerance  $\eta$ 
6 :     Set  $k = k + 1$ 
7 :   else
8 :     Set  $k = k - 1$ 
9 :   end
10: until  $k \leq M$ 
```

# Numerical treatment

FETI formulation (FARHAT & ROUX, 1991; KRUIS & BITTNAR, 2008; DOSTÁL; 2009)



- Lagrange function

$$L(\mathbf{u}, [\mathbf{u}], \lambda) = \frac{1}{2} \mathbf{u}^T \mathbf{K} \mathbf{u} + \frac{1}{2} [\mathbf{u}]^T \mathbf{k} [\mathbf{u}] - \mathbf{d}^T \mathbf{u} + \lambda^T (\mathbf{B}_e \mathbf{u} - \mathbf{u})$$

- Saddle-point problem

$$(\mathbf{u}, [\mathbf{u}], \lambda) = \arg \max_{\lambda} \min_{\mathbf{u}} \min_{\mathbf{B}_i [\mathbf{u}] \geq 0} L(\mathbf{u}, [\mathbf{u}], \lambda)$$

- The “unresolved” box constraint  $\mathbf{B}_i [\mathbf{u}] \geq 0$  will be enforced by a heuristic active set strategy

# Numerical treatment

Dual problem (FARHAT & Roux, 1991; KRUIS & BITTNAR, 2008; DOSTÁL; 2009)

- Optimality conditions

$$[\![\mathbf{u}]\!] = \mathbf{k}^{-1} \boldsymbol{\lambda} \quad \mathbf{u} = \mathbf{K}^+ (\mathbf{d} - \mathbf{B}_e^T \boldsymbol{\lambda}) + \mathbf{R}\boldsymbol{\alpha}$$

$$\mathbf{R}^T (\mathbf{d} - \mathbf{B}_e^T \boldsymbol{\lambda}) = 0 \quad \mathbf{B}_e \mathbf{u} - [\![\mathbf{u}]\!] = 0$$

- Eliminate primary variables  $\mathbf{u}$  and  $[\![\mathbf{u}]\!]$  ( $\mathbf{R}^T \mathbf{d} = 0$ )

$$\begin{bmatrix} \mathbf{B}_e \mathbf{K}^+ \mathbf{B}_e^T + \mathbf{k}^{-1} & -\mathbf{B}_e \mathbf{R} \\ -\mathbf{R}^T \mathbf{B}_e^T & 0 \end{bmatrix} \begin{bmatrix} \boldsymbol{\lambda} \\ \boldsymbol{\alpha} \end{bmatrix} = \begin{bmatrix} \mathbf{B}_e \mathbf{K}^+ \mathbf{d} \\ 0 \end{bmatrix}$$

- Problem reduced to **interfaces** ✓
- Efficiently** solvable using Projected Conjugate Gradient solvers
- Perfect bonding  $\mathbf{k}^{-1} \rightarrow 0$  ✓

# Numerical treatment

Active set strategy (GRUBER & ZEMAN, in preparation)



## 1 Contact detection algorithm

- Tag interfacial nodes as primal ( $\mathcal{P}$ ) or dual ( $\mathcal{D}$ )

$$[\![\mathbf{u}]\!]_e = k_e^{-1} \lambda_e$$

- For  $\mathcal{D}$  : determine contact nodes ( $\mathcal{C}$ ) from  $\lambda$
- For  $\mathcal{P}$  : determine contact nodes ( $\mathcal{C}$ ) from  $[\![\mathbf{u}]\!]$
- For  $\mathcal{C}$  : enforce  $[\![\mathbf{u}]\!]_n = 0$  by setting compliance to zero
- For  $\mathcal{P} - \mathcal{C}$  : enforce  $\lambda = 0$  by static condensation

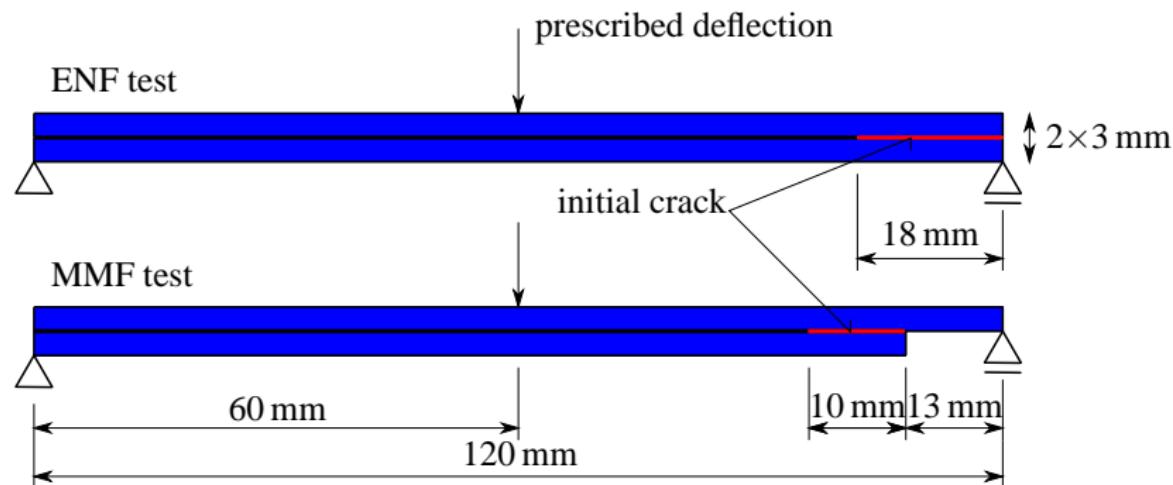
## 2 Solve dual system by standard Projected Conjugate Gradients

## 3 Check convergence of Lagrange multipliers $\lambda$

- Usually converges in 3–4 iterations ✓

# Examples

Notched flexure tests (VALOROSO & CHAMPANEY, 2006)



<i>Mesh size h</i>	<i># of primal DOFs</i>	<i># of dual DOFs</i>
1 mm	1,936	242 (12%)
0.75 mm	3,220	322 (10%)
0.5 mm	6,748	482 (7%)

- $M = 40$  time steps

# Examples

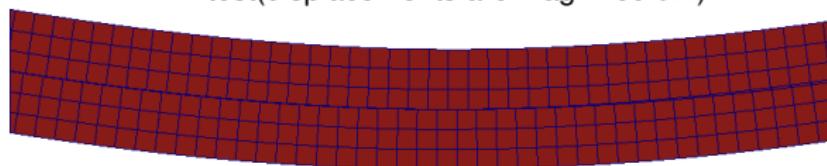
## Material data

- Domains:  $E = 75 \text{ GPa}$ ,  $\nu = 0.3$
- Interfaces: Piecewise linear interfacial law

<i>Material parameter</i>	<i>Ductile</i>	<i>Brittle</i>
$k_n = k_t (\text{GNm}^{-3})$	125	$125 \times 10^6$
$G_n^o = G_t^o (\text{Jm}^{-2})$	100	0.01
$G_n^c = G_t^c (\text{Jm}^{-2})$	250	25
Interaction parameters $a_i, b_i$	2	2

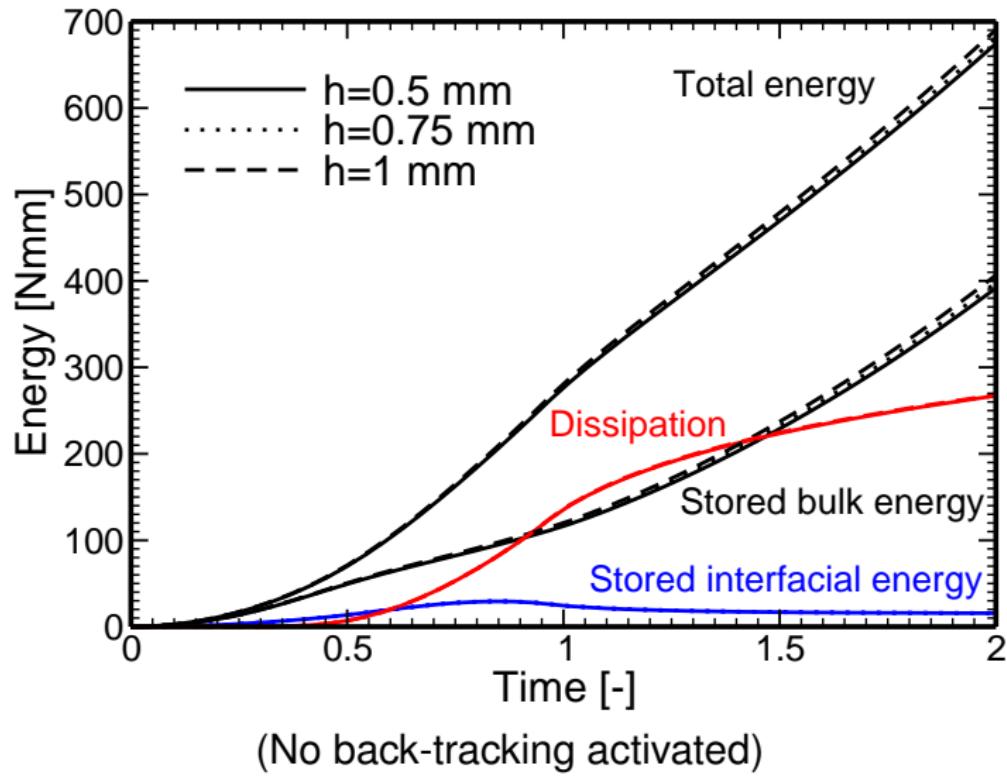
- Algorithmic settings:  $\delta = 10^{-5}$ ,  $\eta = 10^{-3}$
- In-house experimental MATLAB implementation

ENF test(displacements are magnified 5×)



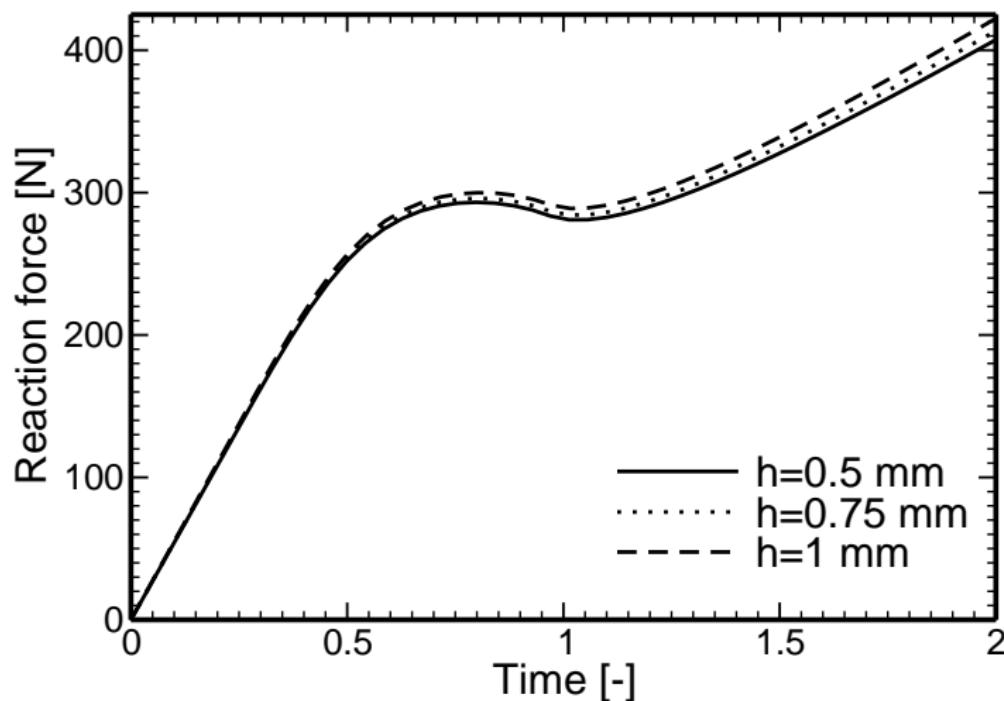
# Examples

Energetics of ENF test,  $h \rightarrow 0$ , ductile interface, 60–300 s



# Examples

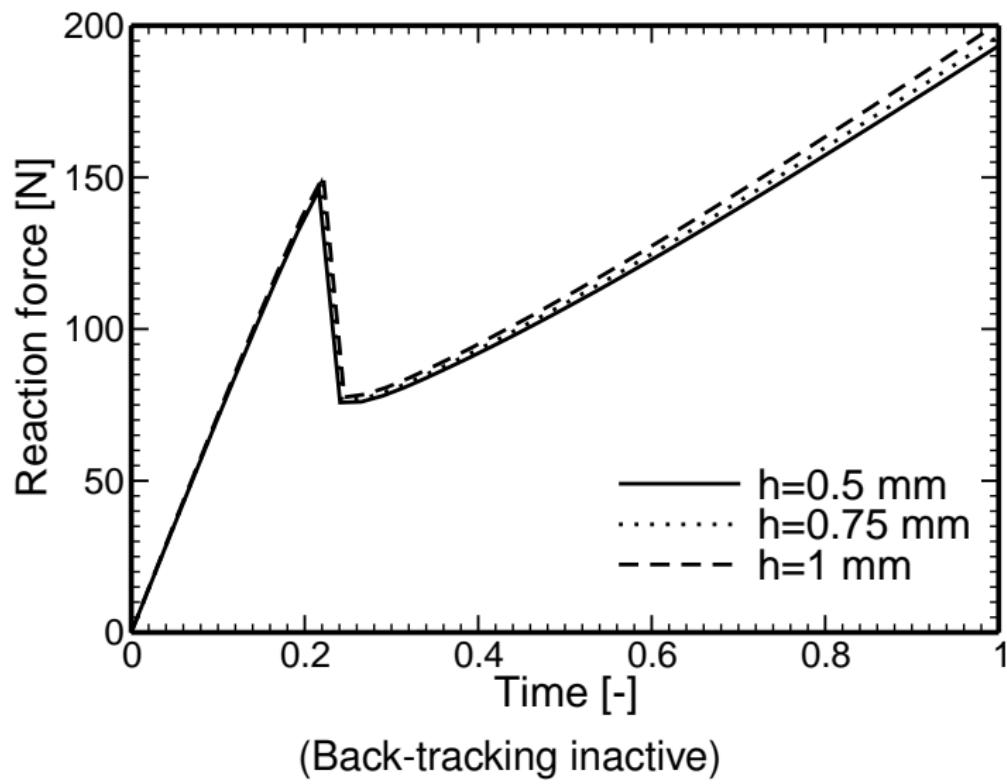
Force-displacement curves of ENF test,  $h \rightarrow 0$ , ductile interface, 60–300 s



(No back-tracking activated)

# Examples

Force-displacement curves of ENF test,  $h \rightarrow 0$ , brittle interface, 60–300 s



# Examples

ENF test, ductile interface, snapshots of delamination evolution



$t = 0.33$



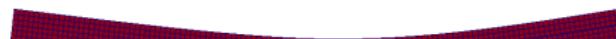
$t = 0.50$



$t = 0.66$



$t = 1.00$



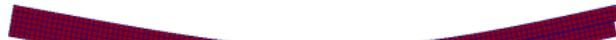
$t = 1.17$



$t = 1.33$



$t = 1.50$



$t = 1.67$



$t = 1.83$

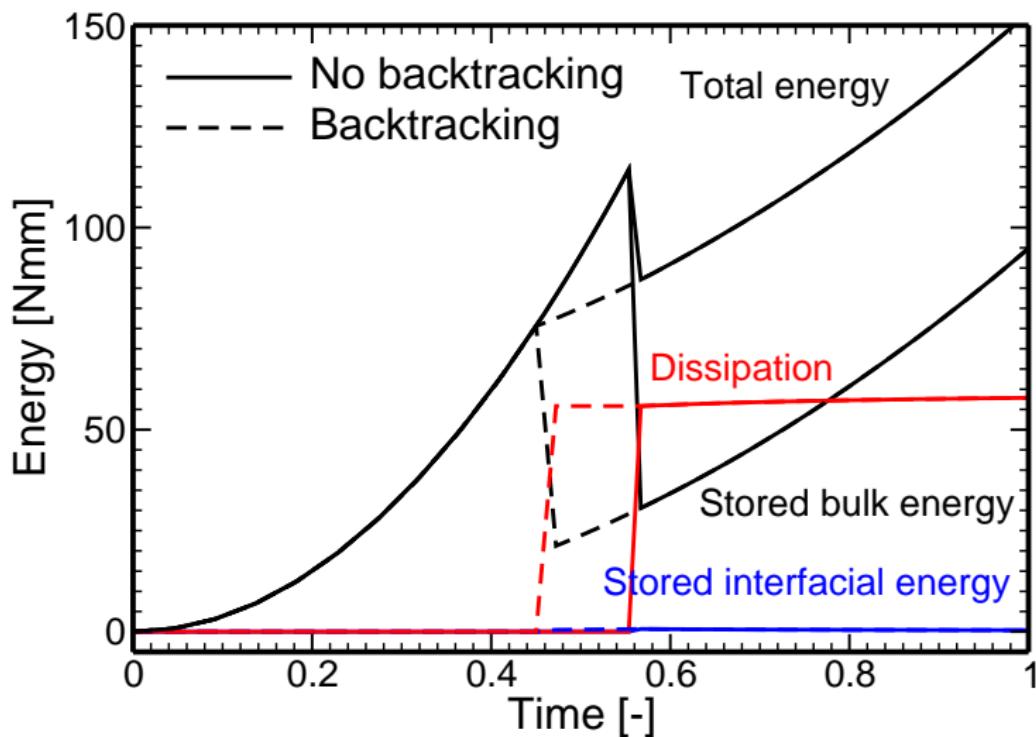


$t = 2$

Displacement are scaled 5×

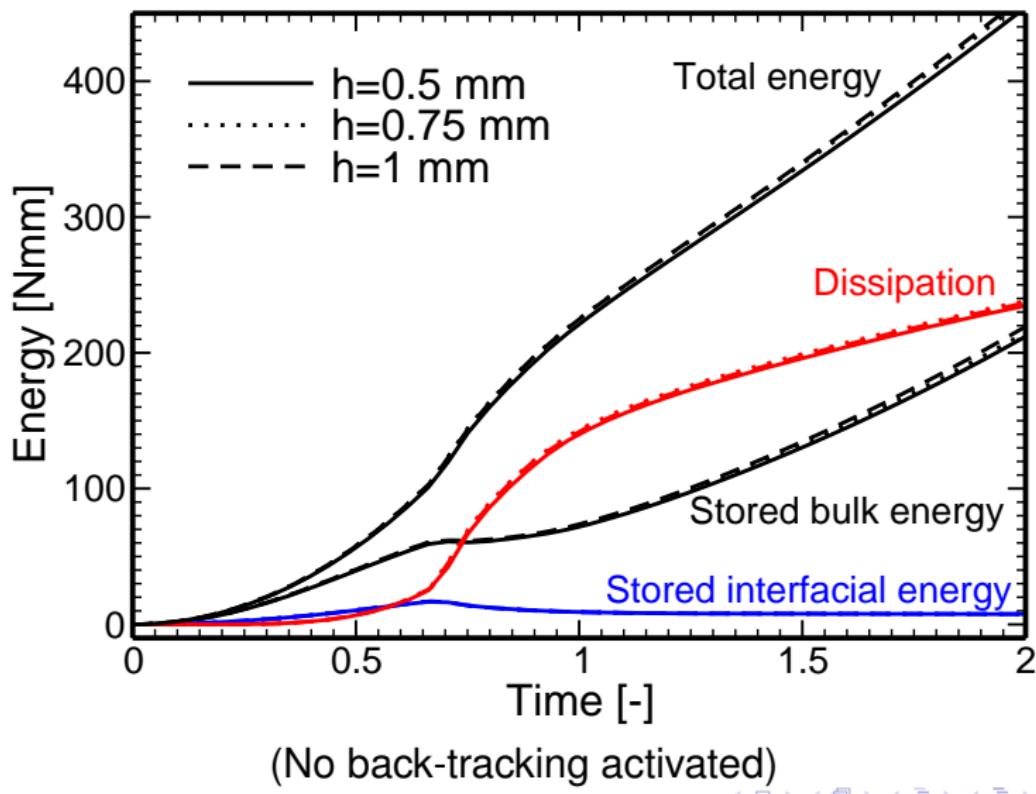
# Examples

Energetics of modified ENF test,  $h = 0.1$  mm, brittle interface



# Examples

Energetics of MMF test,  $h \rightarrow 0$ , ductile interface



# Examples

## Snapshots of MMF test



$t = 0.27$



$t = 0.44$



$t = 0.62$



$t = 0.96$



$t = 1.14$



$t = 1.31$



$t = 1.48$



$t = 1.65$



$t = 1.83$



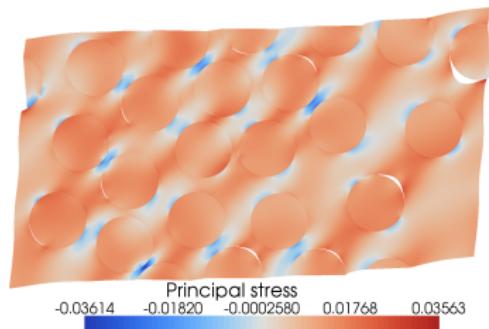
$t = 2$

Displacement are scaled 5×

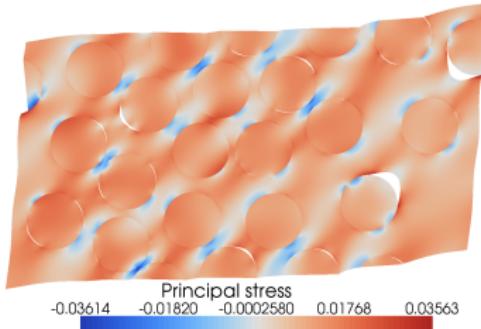
# Extensions

## Debonding in fibre-reinforced composites

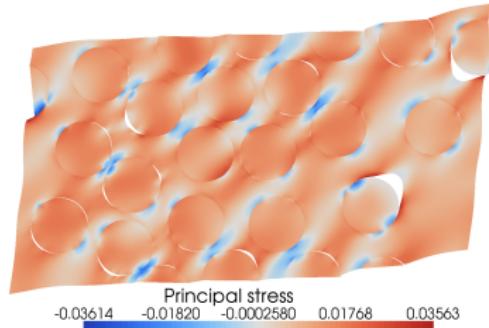
$t = 0.4$



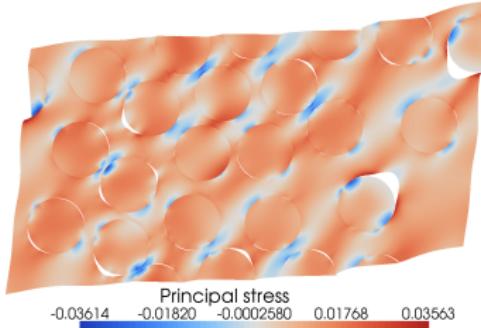
$t = 0.6$



$t = 0.8$

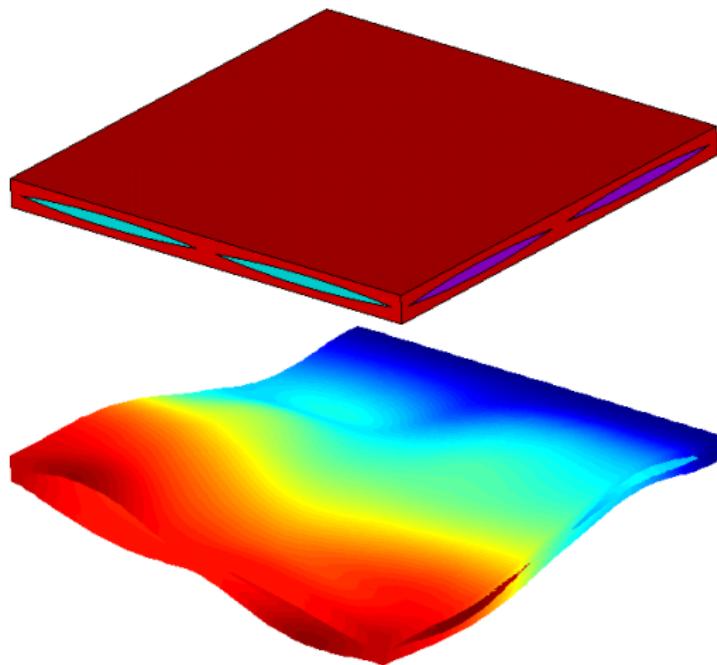


$t = 1$



# Extensions

Woven composites, preliminary results (courtesy of T. KOZUBEK, TU Ostrava)



- Analyzed using MATSOL library

- Energetic framework
  - is well-suited for analysis of inelastic solid mechanics problems
  - provides background for efficient numerical implementation
  - addresses the problems of *stability* of solution when the response is non-unique
- FETI-based algorithm is well-suited for analysis in composite materials, even in sequential mode
- Future extensions
  - Abstract approximation results, homogenization
  - More efficient duality solvers
  - Application to woven composite materials

## Acknowledgements

- Tomáš Roubíček and Alexander Mielke
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