Dva přístupy k vynucení teploty při chlazení proudící vodou.

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Technical motivation

Shape optimization of the cooling plunger cavity

Technical motivation









We denote

$$F_2^e(x) = \begin{cases} 0 & \text{for } x \in [0, x_2^e] \\ f_2^e(x) & \text{for } x \in [x_2^e, 1] \end{cases} ,$$
 (1)

where $x_2^e \in [s_{\min}, 1]$, $(s_{\min} > 0$ is fixed constant given by the minimal thickness of the plunger wall), $f_2^e \in C^{(0),1}([x_2^e, 1])$, $f_2^e(x_2^e) = 0$ and $0 < f_2^e(x) < f_1(x) - s_{\min}$, where f_1 is fixed given increasing function, which represents outward shape of the plunger. Further we assume that $a < f_2^e(x) - s_2$ for $x \in [x_3^e, 1]$, where a > 0 represents radius of supply tube and $s_2 > 0$ minimal admissible split width between the inner wall of plunger cavity and the supply tube, $x_3^e \in [x_2, 1]$ deepness of insertion of the tube.



We denote the set of admissible functions as

$$\begin{split} U_{ad}^{e} &= \left\{ \begin{array}{ll} F_{2}^{e}(x) \in C^{(0),1)}([0,1]) \, ; \ F_{2}^{e}(x) = \left\{ \begin{array}{ll} 0 & \text{for} \quad x \in [0, \, x_{2}^{e}] \\ f_{2}^{e}(x) & \text{for} \quad x \in [x_{2}^{e}, \, 1] \end{array} \right. , \\ & x_{2}^{e} \in [s_{\min}, \, 1], \ s_{\min} > 0, \ f_{2}^{e} \in C^{(0),1}([x_{2}^{e}, \, 1]), \ f_{2}^{e}(x_{2}^{e}) = 0, \\ & 0 < f_{2}^{e}(x) < f_{1}(x) - s_{\min}, \ f_{1} \text{ given}, \\ & a < f_{2}^{e}(x) - s_{2} \text{ for } x \in [x_{3}^{e}, 1], \ a > 0, \ s_{2} > 0, \ x_{3}^{e} \in]x_{2}, 1] \right\}, \end{split}$$

where the function F_2^e describes the technological constraint for inner cavity surface. We assume the region Ω_{Pl}^e , which depends on the design function $F_2^e(x)$, and which is defined by the formula

$$\Omega_{Pl}^e = \{ (x, r) \in \mathbb{R}^2 ; F_2^e(x) < r < f_1(x), \text{ for } x \in [0, 1] \} ,$$

where f_1 is fixed given increasing function, which represents inner shape of plunger.





We denote $\Omega^S = \Omega^e_{Pl} \cup \Omega_{Gl} \cup \Omega^e_{Ca} \cup \Omega_{Mo}$.

In the three dimensional region G_{Ca}^e , which is created by rotation of Ω_{Ca}^e around x axis, we assume an axisymmetric incompressible potential flow of water, which is axisymmetric with x axis.

The potential Φ is given as a solution of Neumann problem

$$\Delta \Phi = 0 \qquad \text{in } G_{Ca}^{e} , \qquad (2)$$

$$\frac{\partial \Phi}{\partial n} = g \qquad \text{on } \partial G_{Ca}^{e} , \qquad (3)$$

where $g \in L^2(\partial G_{Ca}^e)$, representing normal component of velocity at entrance to and exit of plunger cavity, is in the form

$$g = \begin{cases} 0 & \text{on} \quad \Gamma_2^{3D} \cup \Gamma_5^{3D} \\ h_{velo}^{in} & \text{on} \quad \Gamma_{in}^{3D} \\ h_{velo}^{out} & \text{on} \quad \Gamma_{out}^{3D} \\ h_{velo}^{out} & \text{on} \quad \Gamma_{out}^{3D} \\ \end{cases},$$
(4)

 h_{velo}^{in} normal velocity at the entrance Γ_{in}^{3D} , $(h_{velo}^{in} < 0)$ and h_{velo}^{out} normal velocity at the exit Γ_{out}^{3D} . Further we assume

$$\int_{\Gamma_{in}^{3D} \cup \Gamma_{out}^{3D}} g \, dS = 0 \, . \tag{5}$$

The variational formulation for the potential function has form: We look for the function $\Phi \in H^1(G^e_{Ca})$ such that

$$\int_{G_{Ca}^{e}} \frac{\partial \Phi}{\partial x_{i}} \frac{\partial \varphi}{\partial x_{i}} dV = \int_{\Gamma_{in}^{3D} \cup \Gamma_{out}^{3D}} g\varphi \, dS \quad \forall \varphi \in H^{1}(G_{Ca}^{e}) .$$
(6)

Velocity field of flowing water $\boldsymbol{u} = (u_1, u_2, u_3)$ in cavity G_{Ca}^e is given as

$$u = \operatorname{grad}\Phi$$
 . (7)

Theorem *(existence and uniqueness of velocity field)* Under the assumption (5) there exists unique velocity field of the form (7) satisfying the estimate of Euclid norm in the form

$$\||\boldsymbol{u}\|\|_{L^{2}(G_{Ca}^{e})} \leq c \left(\|h_{velo}^{in}\|_{L^{2}(\Gamma_{in}^{3D})} + \|h_{velo}^{out}\|_{L^{2}(\Gamma_{out}^{3D})}\right) .$$
(8)

Proof See [1].

The energy equation for stationary flow \boldsymbol{u} with steady temperature in three dimensions has the form

$$c_v \operatorname{grad} \vartheta . \boldsymbol{u} - \frac{k}{\varrho} \Delta \vartheta = q$$
 (9)

We put u = 0 in G_{Pl}^e , G_{Gl} and G_{Mo} (regions created by rotation of Ω_{Pl}^e , Ω_{Gl} and Ω_{Mo} around x axis) because of there is no flowing liquid inside. Further we consider q = 0 in G_{Pl}^e , G_{Ca}^e and G_{Mo} (there is no heat sources inside). Denote $G = G_{Pl}^e \cup G_{Gl} \cup G_{Ca}^e \cup G_{Mo}$.

We divide the notion for searched function ϑ representing distribution of temperature in the system to the sum of four functions as

$$\vartheta = \vartheta_0 + \vartheta_1 + \vartheta_2 + \vartheta_3 \; ,$$

where

$$\vartheta_i = \begin{cases} \vartheta|_{G_i} & \text{in } G_i \\ 0 & \text{in } G \setminus G_i \end{cases} \quad \text{for } i = 0, 1, 2, 3 , \qquad (10)$$

 $(G_0 \equiv G_{Pl}^e, G_1 \equiv G_{Gl}, G_2 \equiv G_{Ca}^e, G_3 \equiv G_{Mo}).$ Further we denote by $\vartheta_i|_{\Gamma_j^{3D}}$ the trace of solution ϑ_i on the boundary Γ_j^{3D} for i, j if Γ_j^{3D} is a boundary of G_i .

We assume the following **boundary conditions:** In the entrance the cooling water has constant temperature $15^{\circ}C$

$$\vartheta_2|_{\Gamma_{in}^{3D}} = 288 \quad \text{on } \Gamma_{in}^{3D}.$$

 $\vartheta_2|_{\Gamma_{out}^{3D}} = h_{out}^e \quad \text{on } \Gamma_{out}^{3D}.$

Tube is isolated, thus

$$\frac{\partial \vartheta_2}{\partial n} = 0 \quad \text{on} \ \Gamma_5^{3D} .$$

The boundary condition on Γ_3^{3D} is given as

$$\vartheta_i|_{\Gamma_3^{3D}} = h_3 \quad \text{on} \quad \Gamma_3^{3D} \quad i = 0, 1, 3,$$

Heat-transfer through Γ_2^{3D} (i. e. between plunger and water) is modeled as boundary condition for contact of two bodies

$$\left(-k_0 \frac{\partial \vartheta_0}{\partial n}\right)^- = \left(-k_2 \frac{\partial \vartheta_2}{\partial n}\right)^+ \qquad \text{on} \quad \Gamma_2^{3D} , \qquad (11)$$

The boundary Γ_7^{3D} (i. e. between mould and environment) is modeled as boundary condition of the third kind for contact between body and environment

$$\left(-k_3\frac{\partial\vartheta_3}{\partial n}\right)^- = \alpha(\vartheta_3|_{\Gamma_7} - \vartheta_4) \qquad \text{on} \quad \Gamma_7^{3D} , \qquad (12)$$

We use the transit condition for contact between two bodies, where one of them changes its state of matter because of the influence of solidification to describe heat-transfer through the boundary Γ_1^{3D} between glass piece and plunger. Thus

$$\left(k_1 \frac{\partial \vartheta_1}{\partial n}\right)^+ - \left(k_0 \frac{\partial \vartheta_0}{\partial n}\right)^- = \beta_1 \quad \text{on} \quad \Gamma_1^{3D} , \quad (13)$$

Analogously we describe heat-transfer through the boundary Γ_6^{3D} between glass and mould.

$$\left(k_1 \frac{\partial \vartheta_1}{\partial n}\right)^+ - \left(k_3 \frac{\partial \vartheta_3}{\partial n}\right)^- = \beta_6 \quad \text{on} \quad \Gamma_6^{3D} . \quad (14)$$

Due to rotational symmetry we transform the problem to cylindrical coordinates and use dimensional reduction to x, r coordinates.

By this way we obtain two dimensional velocity field of flowing water $w^e = (w_1, w_2)$ by relations

$$w_1 = u_1 , \qquad (15)$$

$$w_2 = \sqrt{(u_2)^2 + (u_3)^2} , \qquad (16)$$

where $\boldsymbol{u} = (u_1, u_2, u_3)$ is defined in (7).

Dimensional reduction leads to one more boundary condition in axis of the system Γ_4 , which means that there is no heat flow in the normal direction to axis, thus

$$rac{\partial artheta_i}{\partial n} = 0$$
 on Γ_4 $i=0,\,1,\,2,\,3$.

 (\rightarrow)

To define the state problem based on the variational formulation of energy equation in two dimensions we define the operators

(20)

Source<sub>$$\Omega^{S}(\psi) = \varrho_{1} \int_{\Omega_{Gl}} q\psi r \, d\Omega$$
, (21)
Coeff _{$\Omega^{S}(\psi) = \int_{\Gamma_{1}} \beta_{1} \psi r \, d\Gamma + \int_{\Gamma_{6}} \beta_{6} \psi r \, d\Gamma + \int_{\Gamma_{7}} \alpha \vartheta_{4} \psi r \, d\Gamma$, (22)}</sub>

We denote the weighted Sobolev space $H^1_r(\Omega_i)$ with norm

$$\|v\|_{1,r,\Omega_i} = \left(\int_{\Omega_i} \left[\left(\frac{\partial v}{\partial x}\right)^2 + \left(\frac{\partial v}{\partial r}\right)^2 + v^2 \right] r \, d\Omega \right)^{\frac{1}{2}} \quad i = 0, 1, 2, 3 , \qquad (23)$$

 $(\Omega_0 \equiv \Omega_{Pl}^e, \, \Omega_1 \equiv \Omega_{Gl}, \, \Omega_2 \equiv \Omega_{Ca}^e, \, \Omega_3 \equiv \Omega_{Mo}).$

Further we denote

 $H(\Omega) = \{\vartheta; \ \vartheta \text{ defined in (10)}, \ \vartheta_i \in H^1_r(\Omega_i) \text{ for any } i = 0, 1, 2, 3\}$.

We define the norm in $oldsymbol{H}(\Omega)$ as

$$\|\vartheta\|_{\boldsymbol{H}} = \left(\|\vartheta_0\|_{1,r,\Omega_0}^2 + \|\vartheta_1\|_{1,r,\Omega_1}^2 + \|\vartheta_2\|_{1,r,\Omega_2}^2 + \|\vartheta_3\|_{1,r,\Omega_3}^2\right)^{\frac{1}{2}} .$$
(24)

Theorem The set $H(\Omega)$ with the norm (24) is a Hilbert space.

Denote the set

$$\Omega_H = \Omega^S \cup \Gamma_3 \cup \Gamma_{in} \cup \Gamma_{out}$$

and

$${}^{e}\mathcal{H}^{2D} = \left\{ v \in C^{\infty}(\Omega_{H}); v|_{\Gamma_{3} \cup \Gamma_{in} \cup \Gamma_{out}} = 0 \right\}.$$

Let $H_0(\Omega)$ be closure of the set ${}^e\mathcal{H}^{2D}$ according to the norm $H(\Omega)$. We assume existence of the function $\vartheta_{\Gamma}^e \in H(\Omega)$ such that

$$\vartheta^e_{\Gamma}|_{\Gamma_{in}} = 288 \quad \text{on} \quad \Gamma_{in},$$
(25)

$$\vartheta_{\Gamma}^{e}|_{\Gamma_{out}} = h_{out}^{e} \quad \text{on} \quad \Gamma_{out},$$
(26)

$$\vartheta_{\Gamma}^{e}|_{\Gamma_{3}} = h_{3} \quad \text{on} \quad \Gamma_{3},$$
(27)

The State Problem:

We look for the function $\vartheta \equiv \vartheta(F_2^e) \in \boldsymbol{H}(\Omega)$ such that

 $\mathsf{Energy}_{\Omega^S}^{velo}(\vartheta,\,\boldsymbol{w}^e,\,\psi) + \mathsf{Energy}_{\Omega^S}^{cond}(\vartheta,\,\psi) + \mathsf{Environment}_{\Omega^S}(\vartheta,\,\psi) =$

$$= \operatorname{Source}_{\Omega^{S}}(\psi) + \operatorname{Coeff}_{\Omega^{S}}(\psi) \quad \forall \psi \in \boldsymbol{H}_{0}(\Omega) , \qquad (28)$$

$$\vartheta - \vartheta_{\Gamma}^{e} \in \boldsymbol{H}_{0}(\Omega) , \qquad (29)$$

where $F_2^e \in U_{ad}^e$ and w^e is corresponding flow pattern given as gradient of solution (6).

The physical assumption of cooling:

A1 Average temperature of water coming into the plunger cavity is less than average temperature of leaving water.

Theorem The bilinear form (17) satisfies the following condition

$$\operatorname{Energy}_{\Omega^{S}}^{velo}(\vartheta, \, \boldsymbol{w}^{e}, \, \vartheta) > 0 \tag{30}$$

for ϑ , w^e satisfying the physical assumption of cooling A1.

Theorem (existence and uniqueness of solution of the state problem) The state problem (28), (29) has a unique solution $\vartheta(F_2^e)$ for each $F_2^e \in U_{ad}^e$ and associated flow pattern w^e received as gradient of the unique solution of (6) and

$$\|\vartheta(F_2^e)\|_{\boldsymbol{H}} \le \frac{1}{\min(c_0, c_1, c_2, c_3)} \|F_\Omega\|_{\boldsymbol{H}^*} .$$
(31)

Proof See [1].

We will solve the **problem of the optimal design for plunger cavity shape:** We define the **cost functional** as

$$\mathcal{J}^{S}(F_{2}^{e}) = \|\vartheta(F_{2}^{e})|_{\Gamma_{1}} - T_{\Gamma_{1}}\|_{0,r,\Gamma_{1}}^{2}, \qquad (32)$$

where $\vartheta(F_2^e)|_{\Gamma_1}$ is the trace of solution $\vartheta(F_2^e)$ of state problem (28), (29) in the region Ω_{Pl}^e on the boundary Γ_1 , T_{Γ_1} is before chosen fixed constant corresponding to optimal surface plunger temperature. We look for **optimal design** $F_{Opt} \in U_{ad}^e$ such that

$$\mathcal{J}^{S}(F_{Opt}) \leq \mathcal{J}^{S}(F_{2}^{e}) \quad \forall F_{2}^{e} \in U_{ad}^{e} .$$
(33)

Theorem (existence of solution of the problem of optimal design for plunger cavity shape) The optimal design problem (33) has at least one solution.

Proof See [1].

References:

[1] Salač, P.: *Optimal design of the cooling plunger cavity*, Appl. Math. (accepted to publish).

Thank you for your attention.